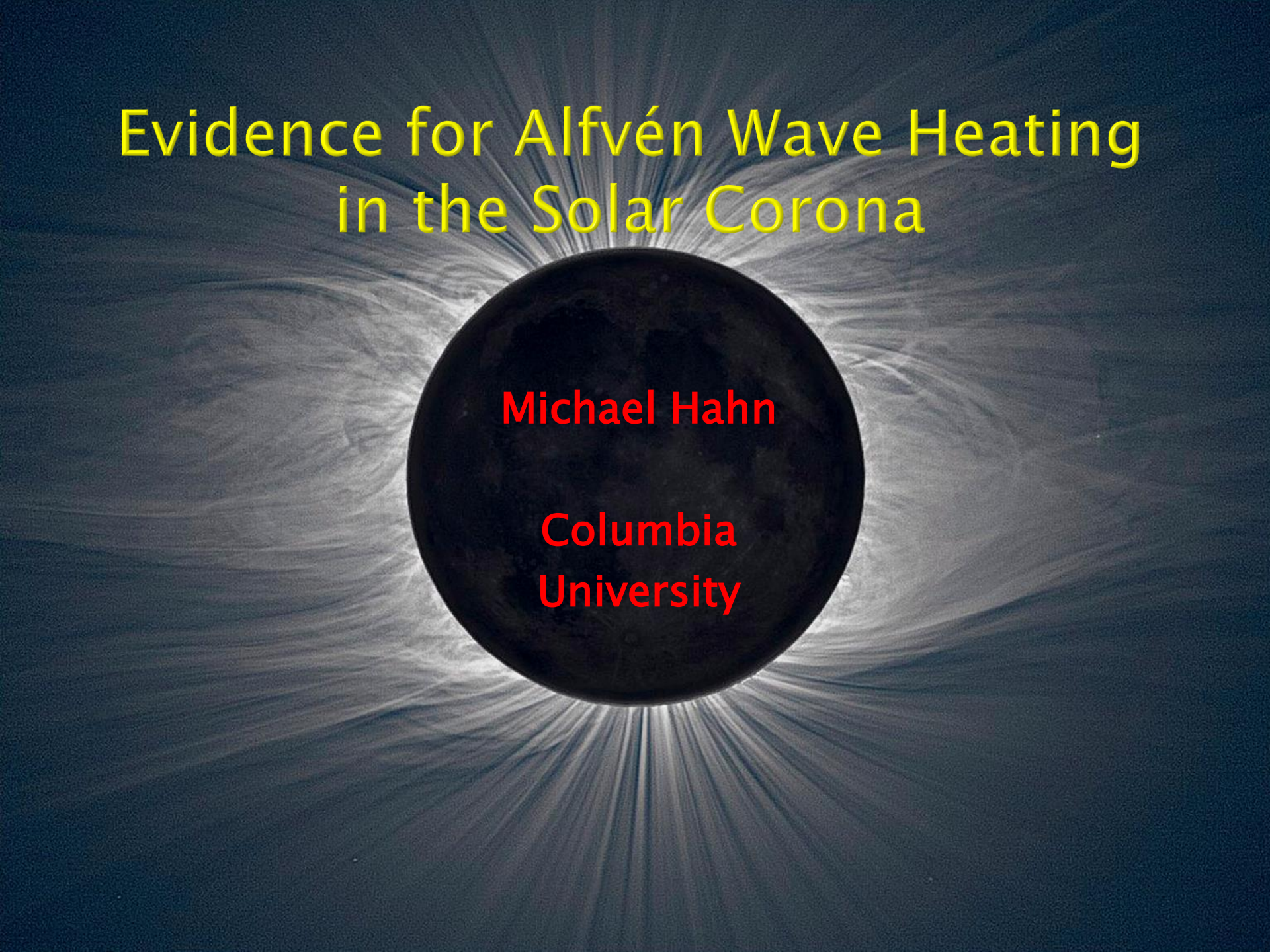


# Evidence for Alfvén Wave Heating in the Solar Corona

**Michael Hahn**

**Columbia  
University**



# Acknowledgements

Daniel Wolf Savin (Columbia)

Enrico Landi (Michigan)

NSF Solar Heliospheric and Interplanetary  
Environment (SHINE)

# Outline

- ▣ Introduction
  - The Coronal Heating Problem.
  - Reconnection.
  - Waves.
- ▣ Spectroscopic measurements of Alfvén waves.
  - Methods – interpreting line widths.
  - Results.
- ▣ Wave damping...
  - due to transverse structure of the magnetic field.
  - due to wave reflection along the field.
- ▣ Conclusions.

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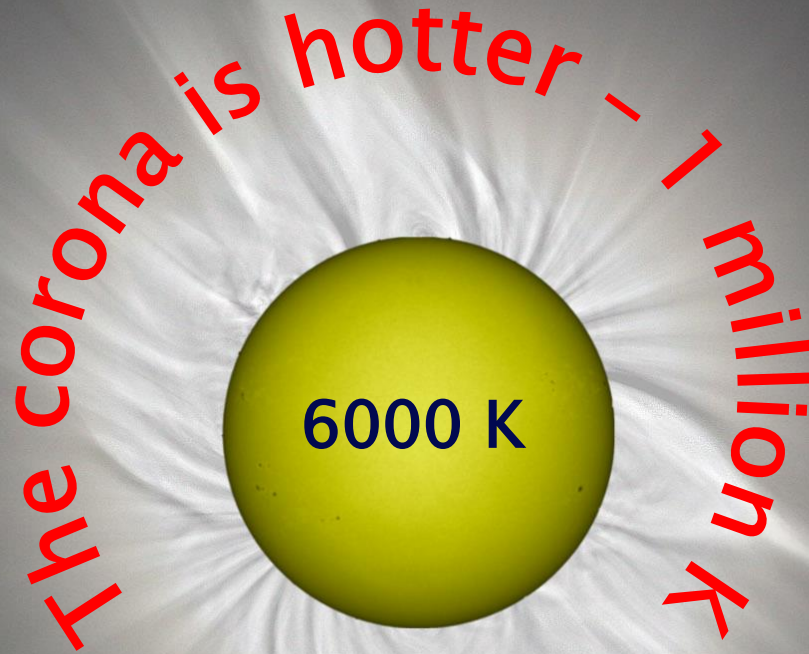


# The Coronal Heating Problem



The Sun is hot – 6000 K

# The Coronal Heating Problem



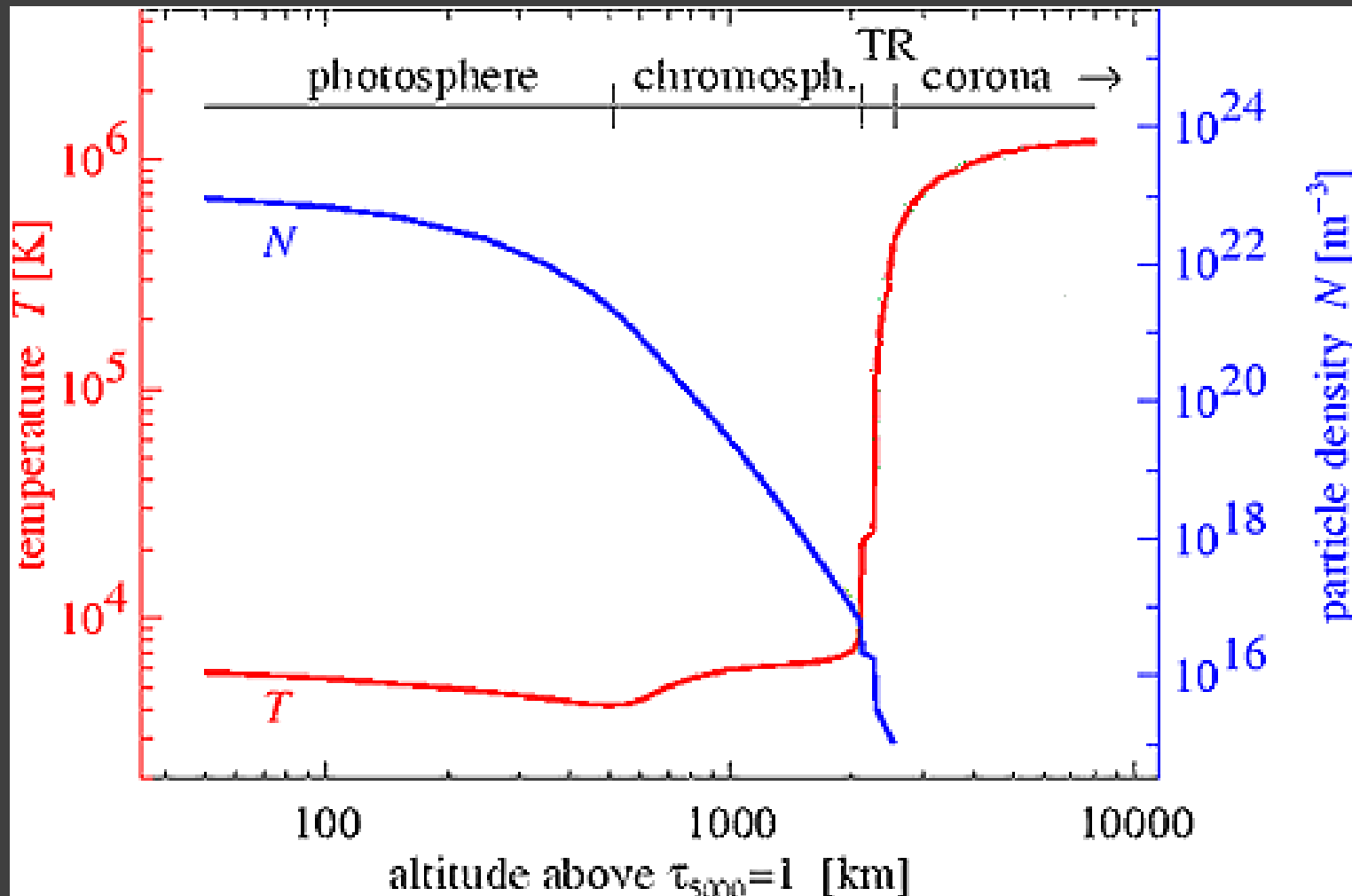
The corona is hotter - 1 million K

6000 K

A diagram of the Sun with a yellow-orange surface. The corona is depicted as a white, wispy, and textured outer layer. The text 'The corona is hotter - 1 million K' is written in red, curved around the top and right sides of the corona. The text '6000 K' is written in blue, centered on the Sun's surface.

# Why?

# The Coronal Heating Problem



Wilhem et al. (2007)

# Driving the Solar Wind



2003 Oct 25 00:00:12



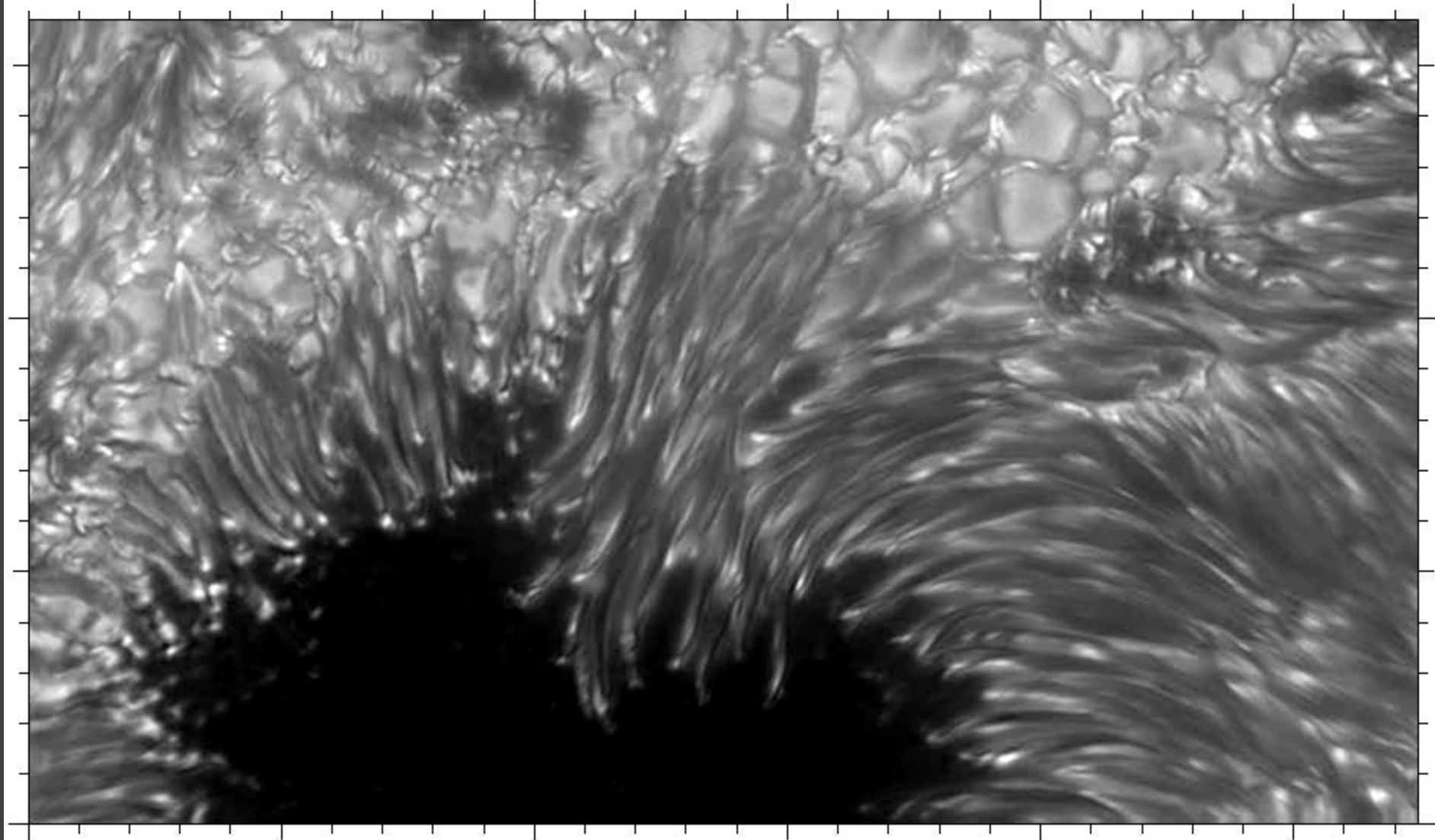
# Coronal Heating

- ▣ Only 0.01% of Sun's energy output required.
- ▣ Energy source is fluid motion in outer convection zone.

# Convection Near a Sunspot

G-band 4305.6+/- 10.8 Å Swedish 1-m Solar Telescope

00:00:00



distance in units of 1000 kilometers

# Coronal Heating

- ▣ Only 0.01% of Sun's energy output required.
- ▣ Energy source is fluid motion in outer convection zone.
- ▣ Jostling of field lines gives energy to corona via:
  - Magnetic reconnection
  - Waves

# So is it Reconnection?

THE ASTROPHYSICAL JOURNAL, 720:824–847, 2010 September 1  
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doi:[10.1088/0004-637X/720/1/824](https://doi.org/10.1088/0004-637X/720/1/824)

## CAN THE SOLAR WIND BE DRIVEN BY MAGNETIC RECONNECTION IN THE SUN'S MAGNETIC CARPET?

STEVEN R. CRANMER AND ADRIAAN A. VAN BALLEGOOIJEN

Harvard-Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, MA 02138, USA; [scanmer@cfa.harvard.edu](mailto:scanmer@cfa.harvard.edu), [avanballegooijen@cfa.harvard.edu](mailto:avanballegooijen@cfa.harvard.edu)

*Received 2010 May 25; accepted 2010 July 13; published 2010 August 13*

### ABSTRACT

The physical processes that heat the solar corona and accelerate the solar wind remain unknown after many years of study. Some have suggested that the wind is driven by waves and turbulence in open magnetic flux tubes, and others have suggested that plasma is injected into the open tubes by magnetic reconnection with closed loops. In order to test the latter idea, we developed Monte Carlo simulations of the photospheric “magnetic carpet” and extrapolated the time-varying coronal field. These models were constructed for a range of different magnetic flux imbalance ratios. Completely balanced models represent quiet regions on the Sun and source regions of slow solar wind streams. Highly imbalanced models represent coronal holes and source regions of fast wind streams. The models agree with observed emergence rates, surface flux densities, and number distributions of magnetic elements. Despite having no imposed supergranular motions in the models, a realistic network of magnetic “funnels” appeared spontaneously. We computed the rate at which closed field lines open up (i.e., recycling times for open flux), and we estimated the energy flux released in reconnection events involving the opening up of closed flux tubes. For quiet regions and mixed-polarity coronal holes, these energy fluxes were found to be much lower than that which is required to accelerate the solar wind. For the most imbalanced coronal holes, the energy fluxes may be large enough to power the solar wind, but the recycling times are far longer than the time it takes the solar wind to accelerate into the low corona. Thus, it is unlikely that either the slow or fast solar wind is driven by reconnection and loop-opening processes in the magnetic carpet.

**Key words:** magnetic fields – magnetohydrodynamics (MHD) – plasmas – solar wind – Sun: corona – Sun: photosphere

# Or Waves?

THE ASTROPHYSICAL JOURNAL, 711:1044–1050, 2010 March 10  
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doi:10.1088/0004-637X/711/2/1044

## DEMONSTRATIONS THAT THE SOLAR WIND IS NOT ACCELERATED BY WAVES OR TURBULENCE

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*Received 2009 July 14; accepted 2010 January 21; published 2010 February 22*

### ABSTRACT

The present work uses observations and theoretical considerations to provide both qualitative and quantitative arguments that hydromagnetic waves, whether turbulent or not, cannot produce the acceleration of the fast solar wind and the related heating of the open solar corona. Waves do exist, and can play a role in the differential heating and acceleration of minor ions, but their amplitudes are not sufficient to power the wind, as demonstrated by extrapolation of magnetic spectra from Helios and *Ulysses* observations. Dissipation mechanisms invoked to circumvent this conclusion cannot be effective for a variety of reasons. In particular, turbulence does not play a strong role in the corona as shown both by observations of coronal striations and other features, and by theoretical considerations of line tying to a nonturbulent photosphere, nonlocality of interactions, and the nature of the kinetic dissipation. We consider possible “ways out” of the arguments presented, and suggest that in the absence of wave or turbulent heating and acceleration, the chromosphere and transition region become the natural source, if yet unproven, of open coronal energization through the production of nonthermal particle distributions.

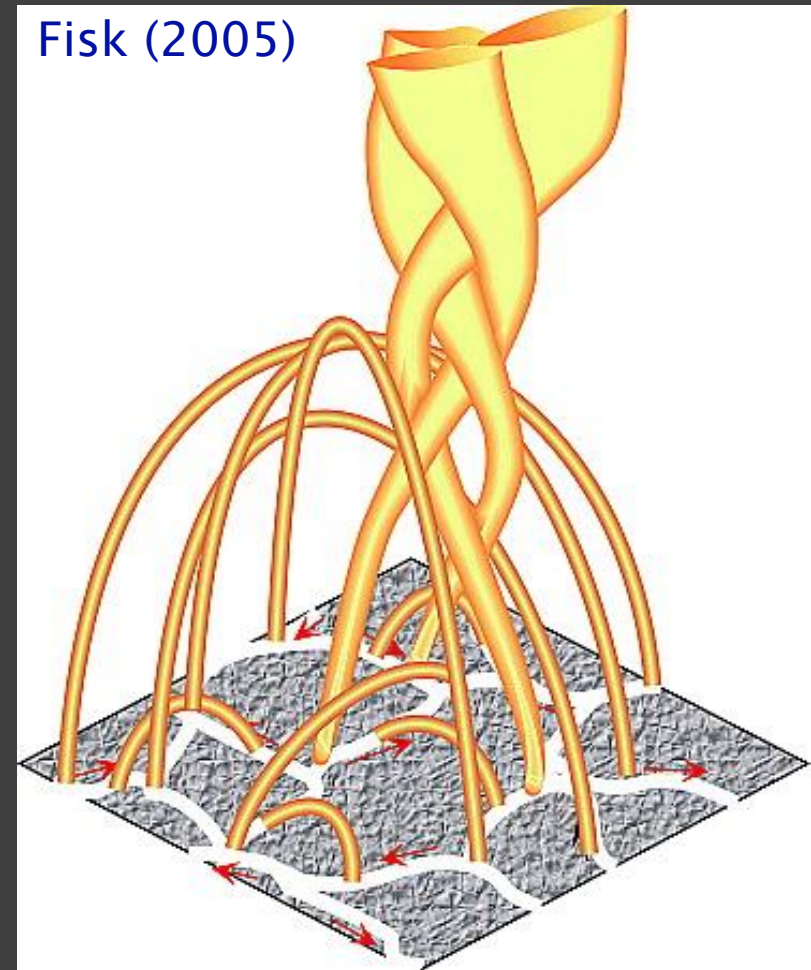
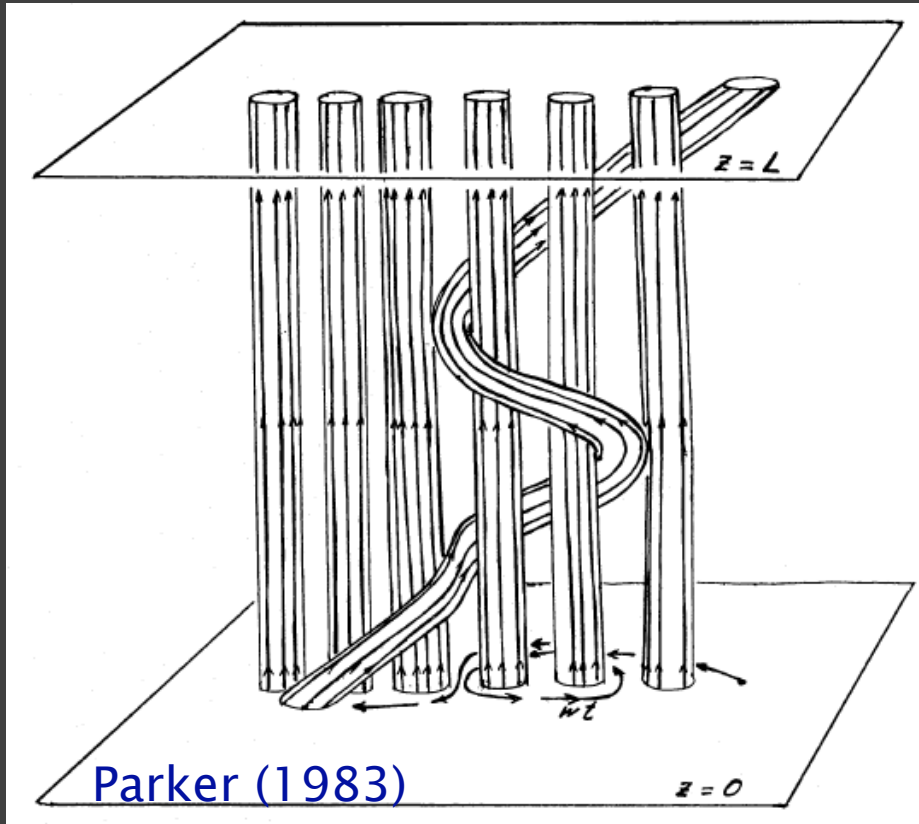
*Key words:* magnetic fields – solar wind – Sun: corona



# Outline

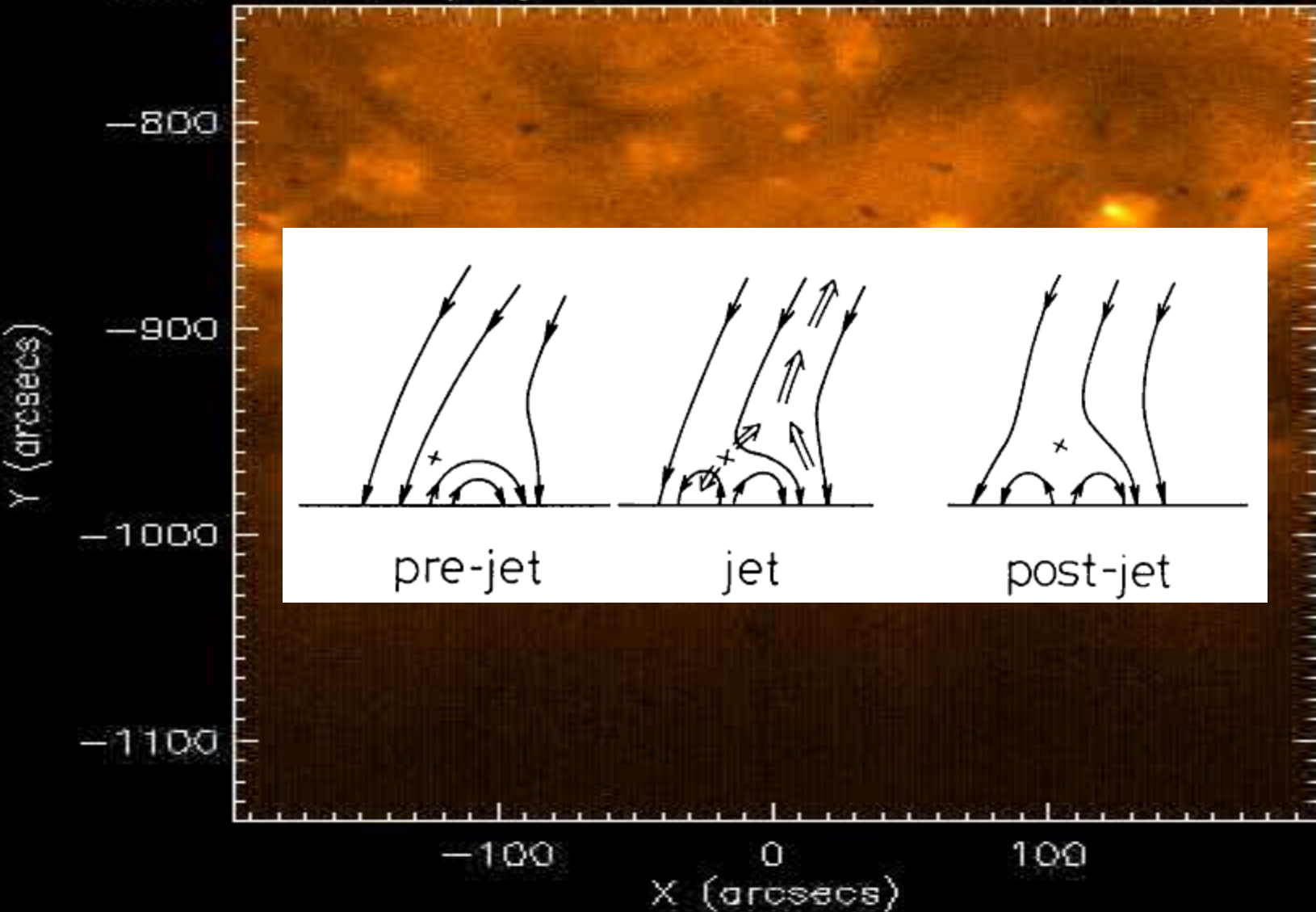
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# Convection tangles magnetic field



# Reconnection Example: X-ray Jets

XRT Thin Al poly 23-Jun-2009 10:14:21.583 UT



# Problems with Reconnection?

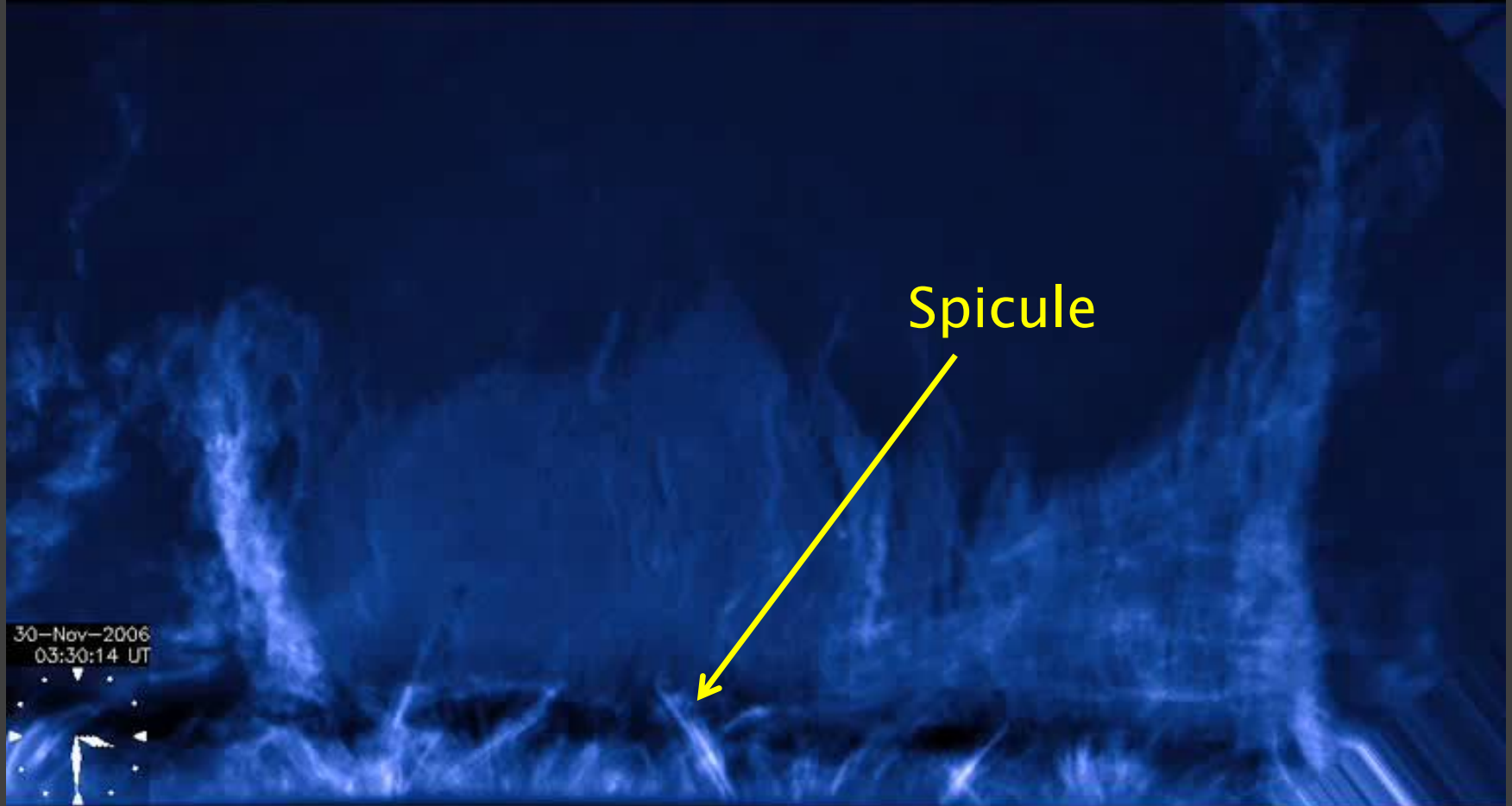
- ▣ Heating looks constant, not episodic.
- ▣ Are there a large enough number of small reconnection events to provide effectively continuous energy input?

# Outline

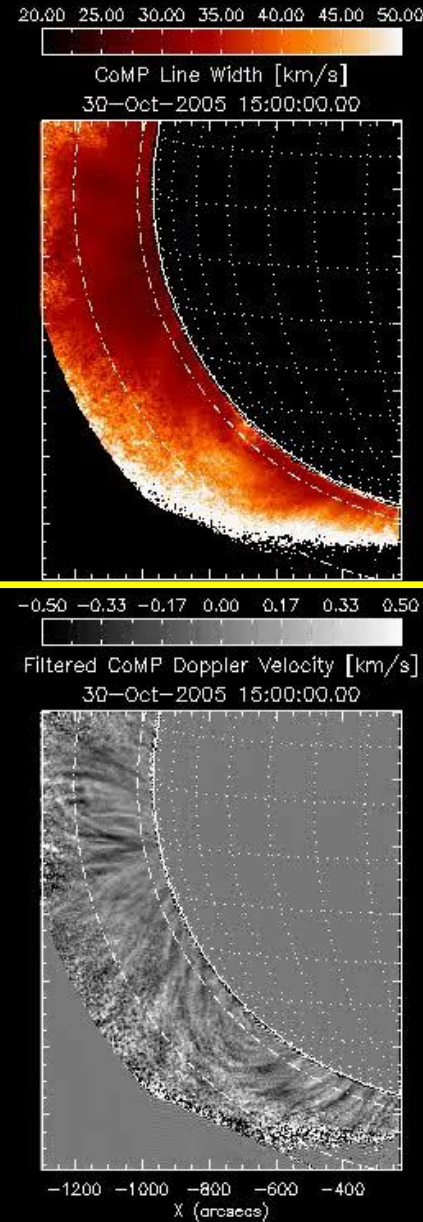
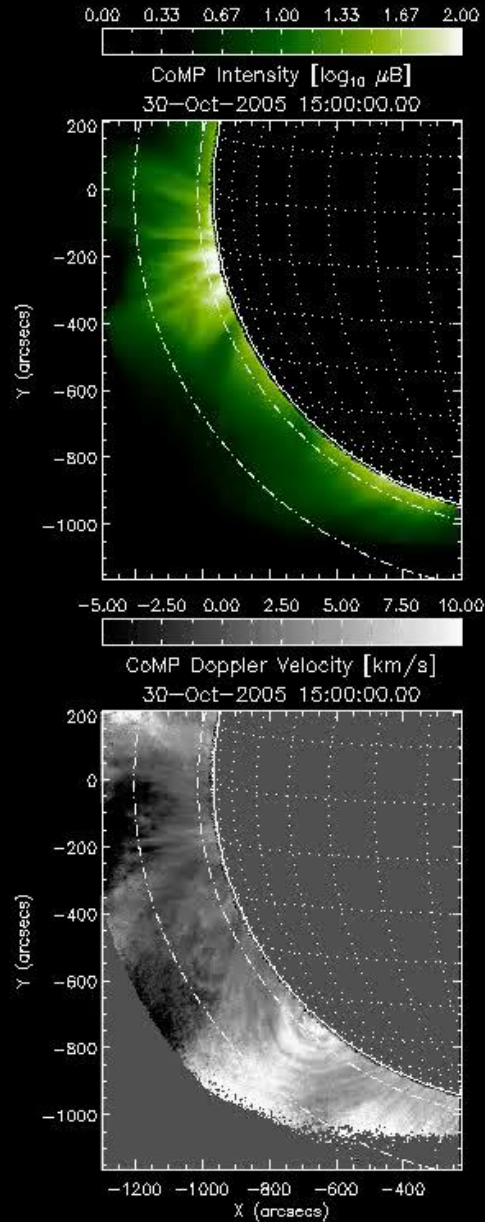
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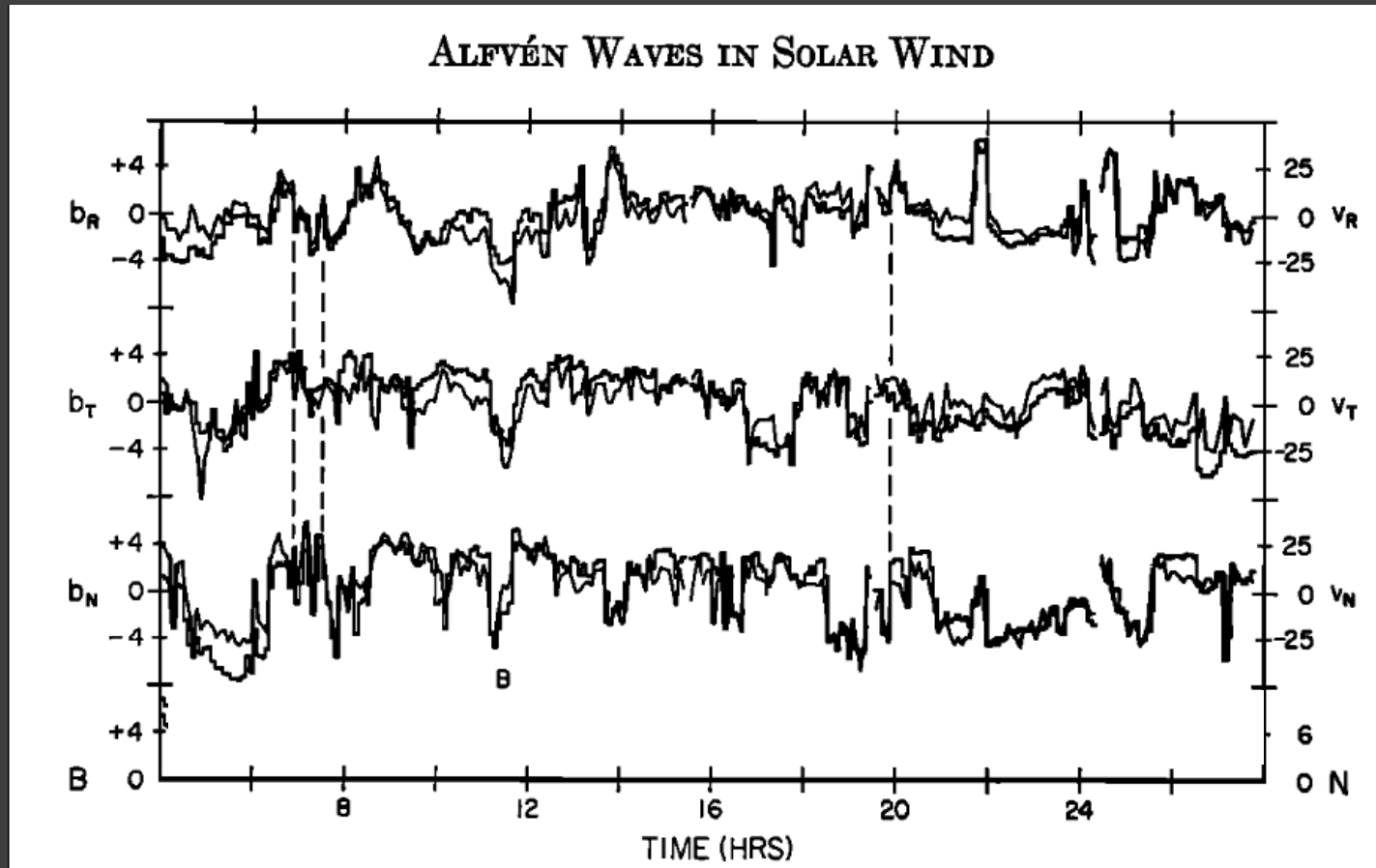
# Alfvén Waves : Chromosphere



# Alfvén Waves : Corona



# Solar Alfvén Waves



Belcher & Davis (1971)

Alfvén waves from Mariner 5 at Venus.

# Wave Heating

- ▣ Waves are found throughout the corona and solar wind.
- ▣ If the waves are damped then they can heat the plasma.
- ▣ Collisions (viscosity, resistivity) damp waves but require length scales of  $\sim 3 R_{\odot}$
- ▣ Above  $\sim 2 R_{\odot}$  low density and supersonic solar wind make heat conduction inefficient.

# Outline

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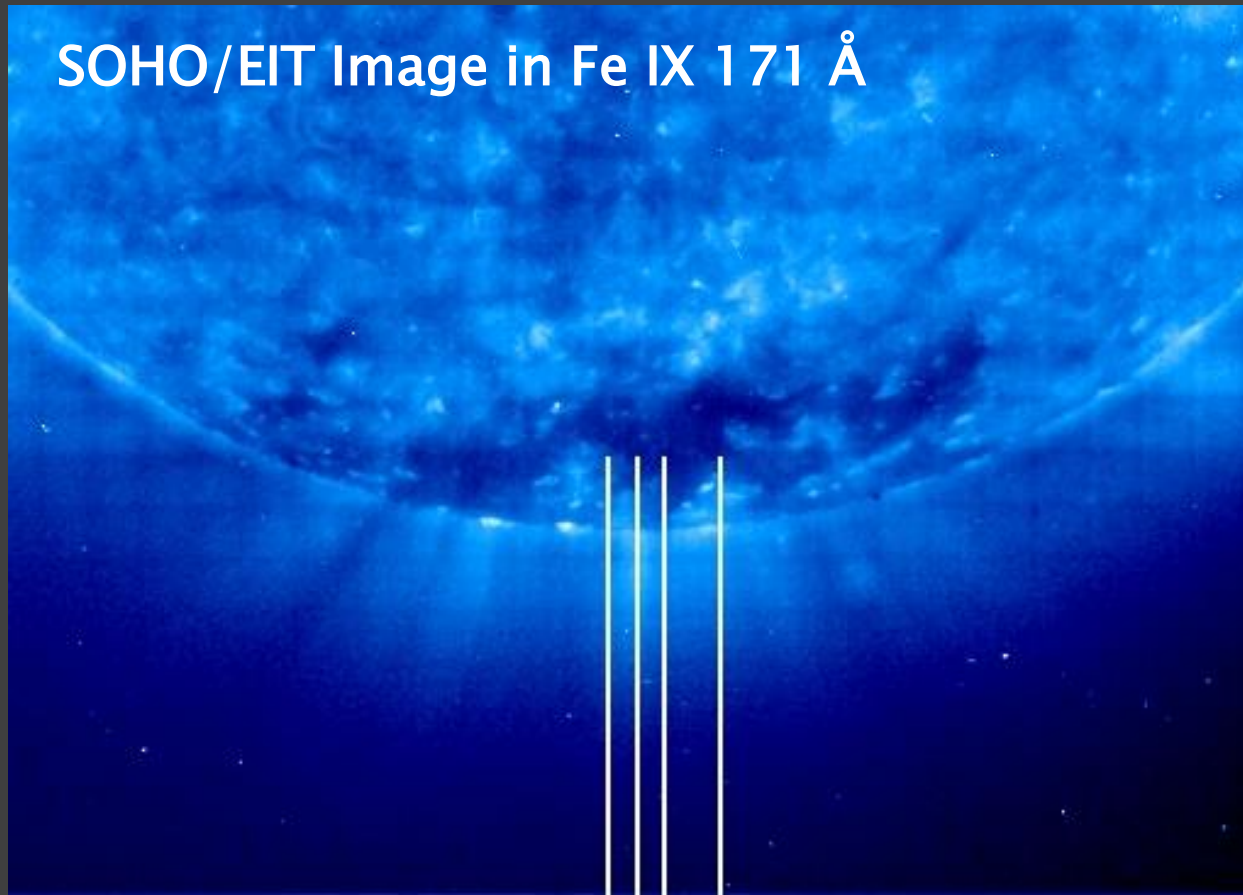


# Extreme Ultraviolet Imaging Spectrometer (EIS)



- ▣ Hinode launched 2006
- ▣ 171–211 Å & 245–291 Å
- ▣ 0.022 Å/pixel
- ▣ 1" and 2" slits

# Observation



- Polar coronal hole from April 2009

# The Solar Corona

Coronal Hole –  
open magnetic  
field lines,  
 $T \sim 1$  MK, less  
dense.

Quiet Sun Corona –  
loop magnetic field  
lines,  $T \sim 1.5$  MK.

Active Region –  
strong magnetic field  
above sunspots,  $T >$   
 $2$  MK.

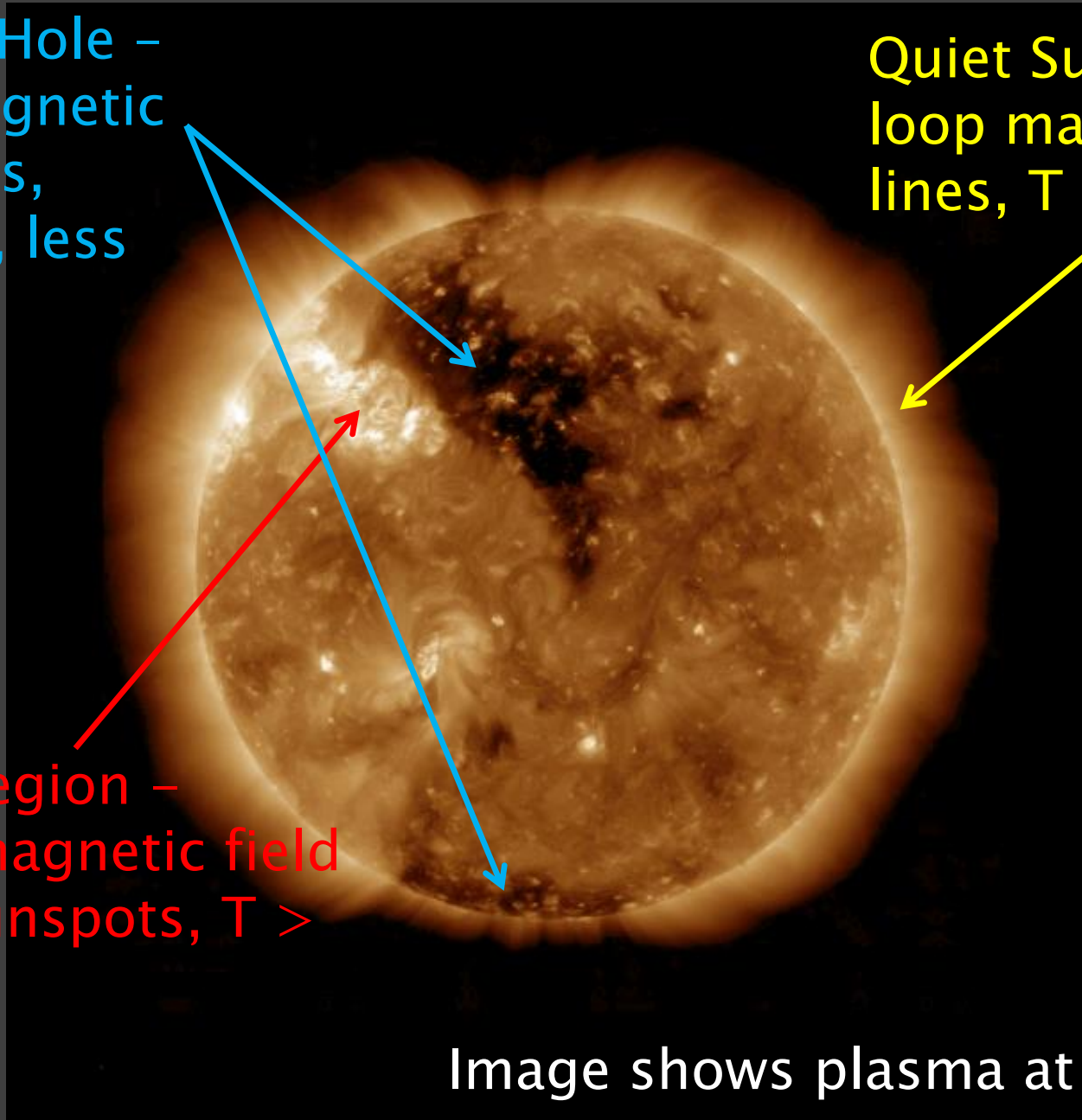
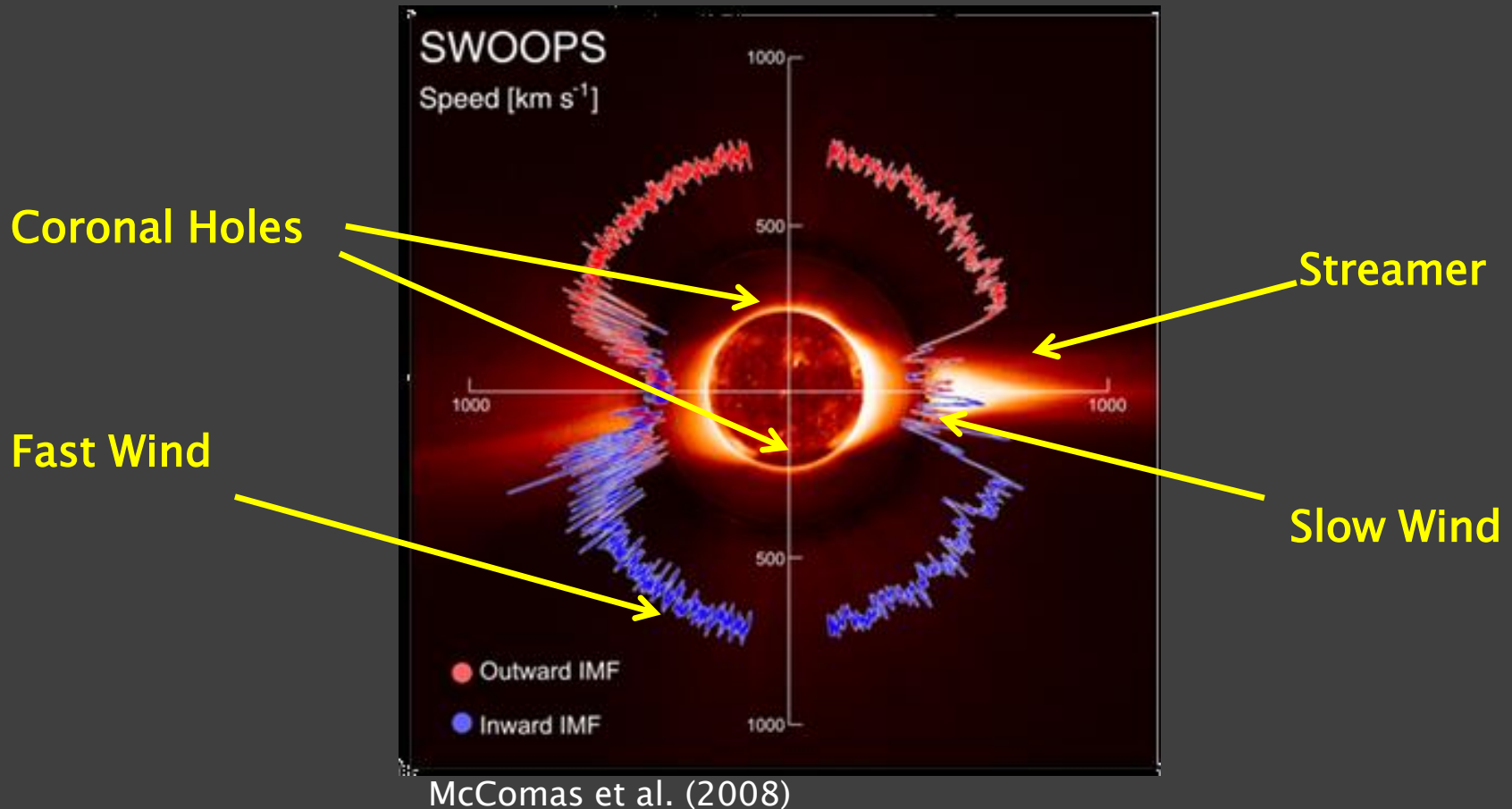


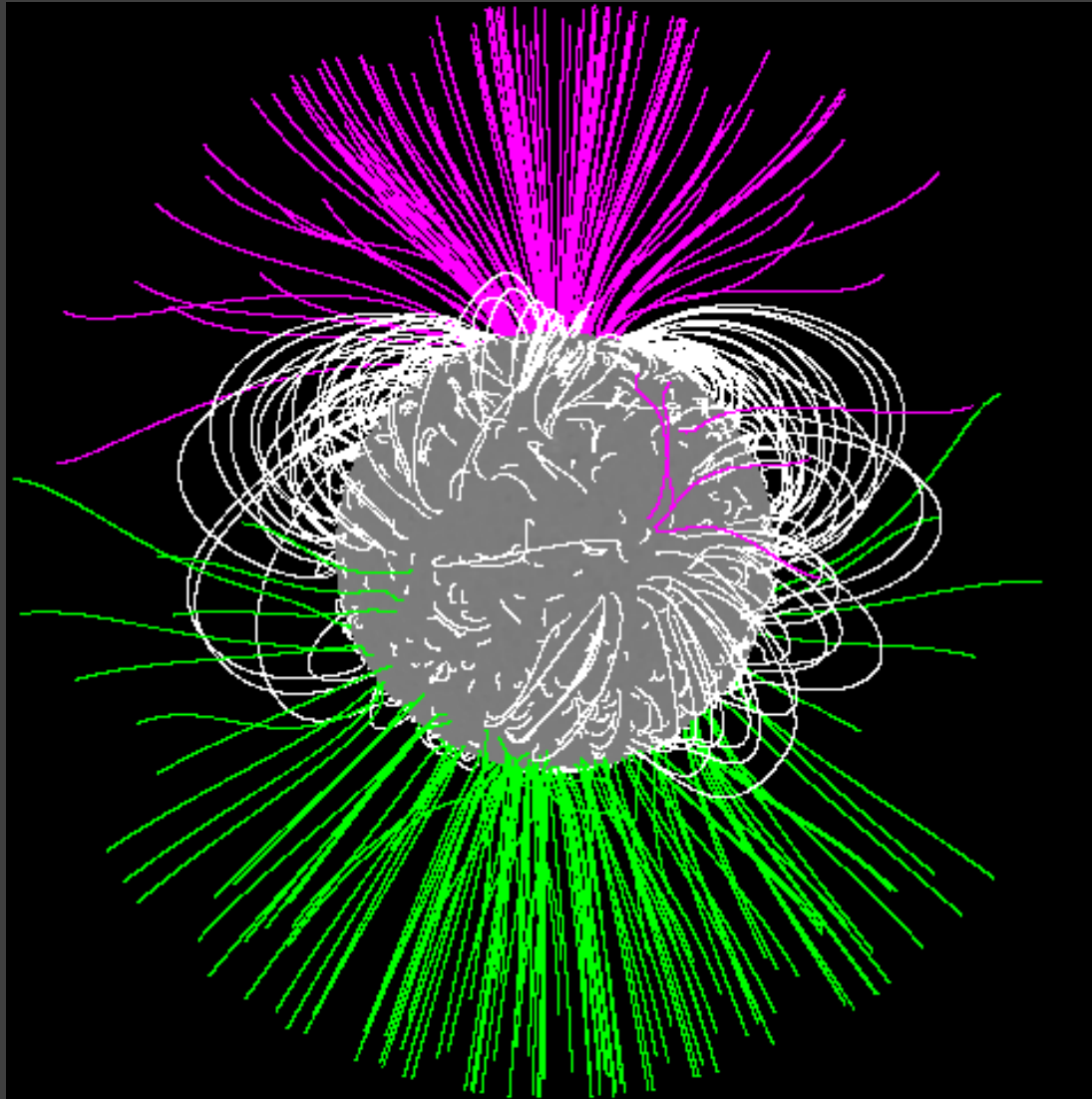
Image shows plasma at  $\sim 1.6$  MK

# Fast solar wind comes from coronal holes



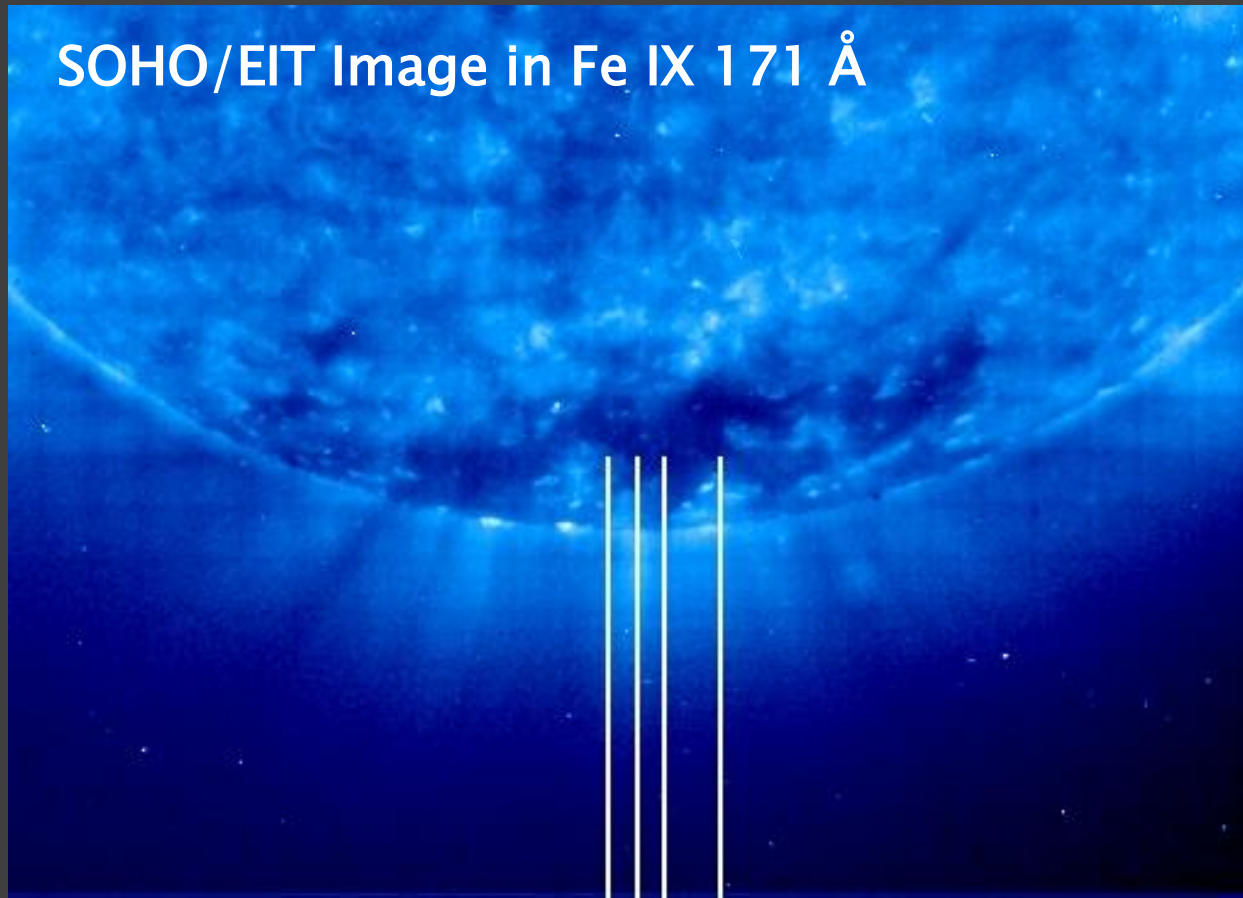
Ulysses – polar orbit of the Sun at ~1 AU.  
SWOOPS = Solar Wind Plasma Experiment

# Solar Magnetic Fields



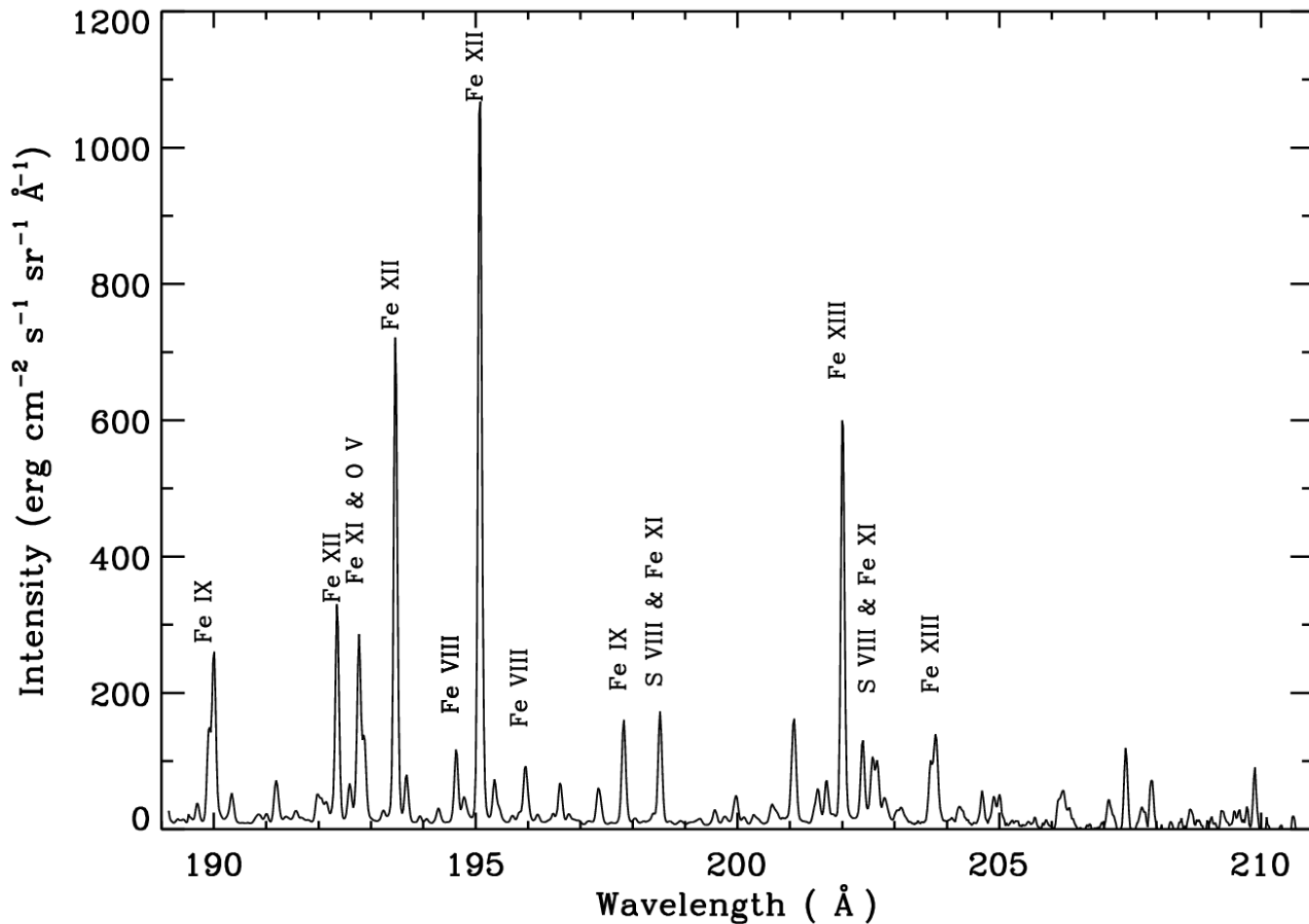


# Observation



- Polar coronal hole from April 2009
- Sum of 4 observations of 30 min. each
- Observed up to  $1.44 R_{\odot}$

# Observation

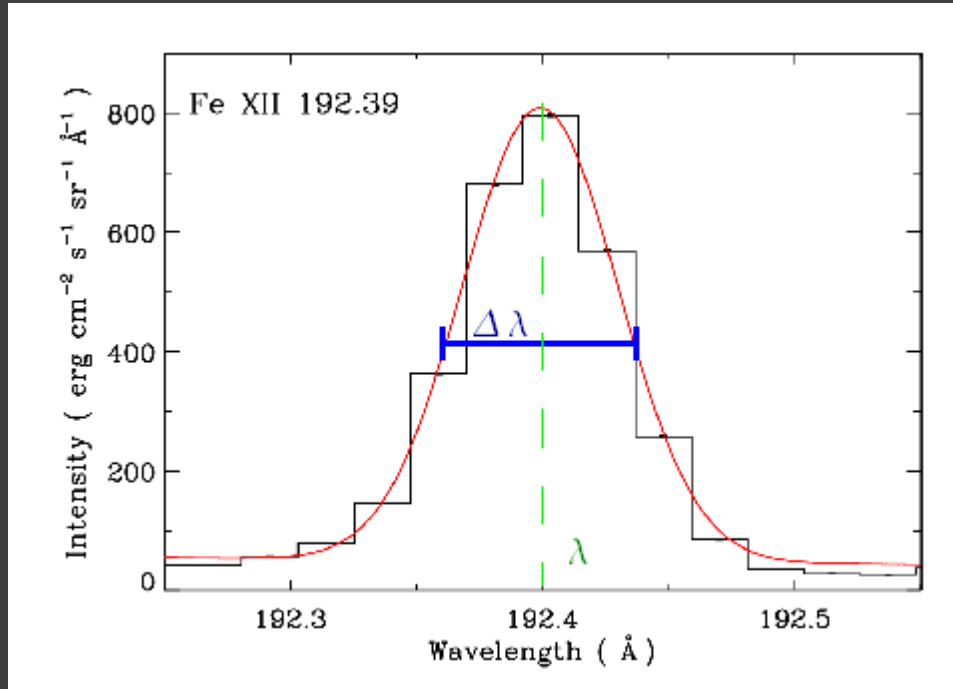


Portion of a typical EIS spectrum.

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# Spectroscopic Signature of Waves



$$\Delta\lambda = \sqrt{\left(\frac{\lambda}{c}\right)^2 \left(\frac{2k_B T_i}{M} + v_{nt}^2\right)}$$

$M$  – Ion Mass

$v_{nt}$  – Nonthermal  
velocity

$T_i$  – Ion temperature

Subtract instrumental width

Wave amplitude  $\delta v$  from  $\langle \delta v^2 \rangle = 2 v_{nt}^2$ .

$T_i$  affected by ion cyclotron resonance heating.

# Energy Conservation

Wave energy flux is

$$F = 2\rho v_{nt}^2 V_A$$

$$V_A = \frac{B}{\sqrt{4\pi\rho}}$$

If undamped then  $FA$  is constant.

$$FA = \frac{1}{\sqrt{\pi}} \rho^{1/2} v_{nt}^2 BA$$

Flux tube  $BA = \text{constant}$ , thus :  $v_{nt} \sim \rho^{-1/4}$

If waves are undamped since  $\rho$  decreases with height  $v_{nt}$  must increase.

# Separate Thermal and Nonthermal

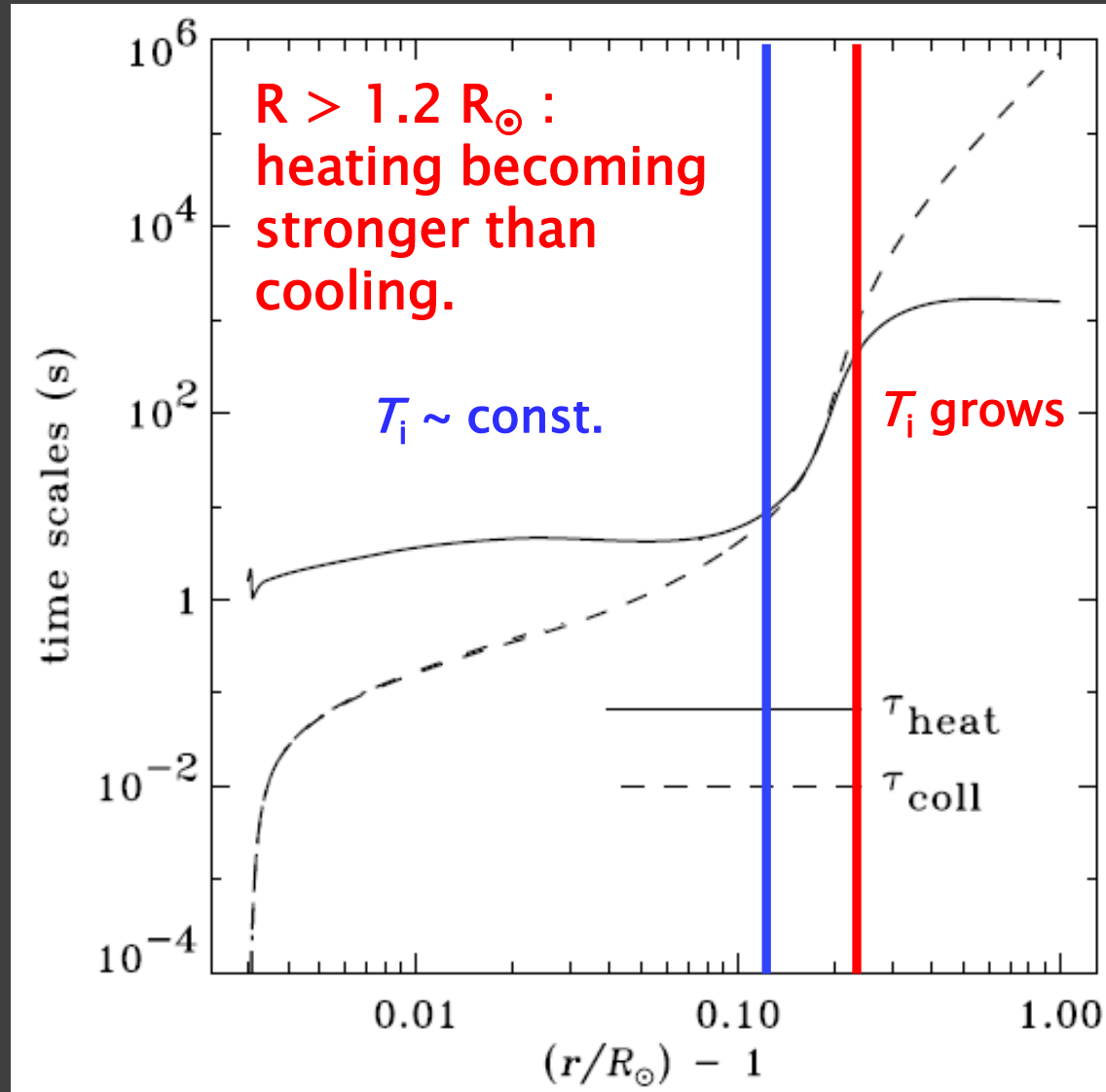
$$v_{\text{eff}}^2 = \sqrt{\frac{2k_{\text{B}}T_{\text{i}}}{M}} + v_{\text{nt}}^2 = \sqrt{v_{\text{th}}^2 + v_{\text{nt}}^2}$$

1. Assume  $T_{\text{i}}$  constant with height.

$$v_{\text{eff}}^2(R_1) - v_{\text{eff}}^2(R_2) = v_{\text{nt}}^2(R_1) - v_{\text{nt}}^2(R_2)$$



# Collision vs. Heating times



Landi & Cranmer (2009)

# Separate Thermal and Nonthermal

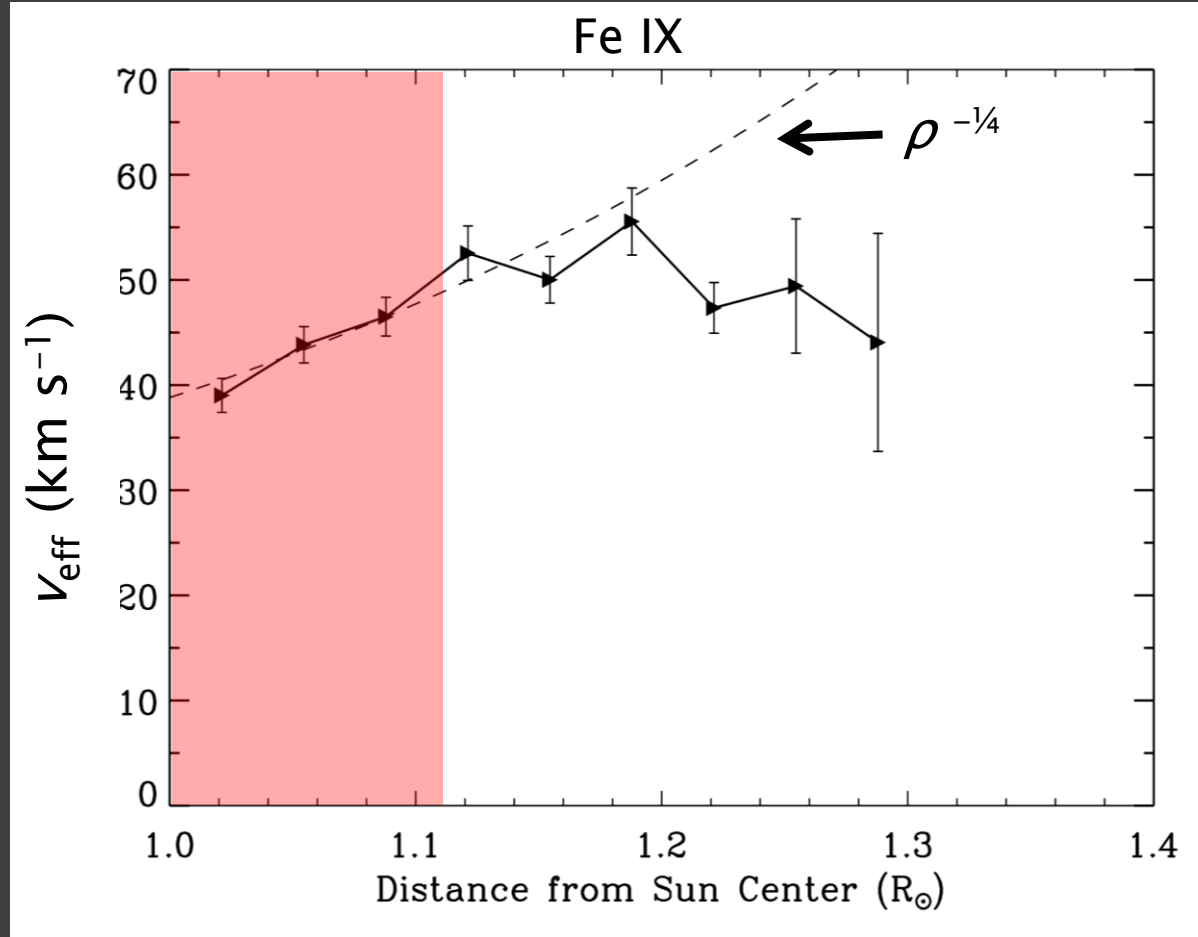
$$v_{\text{eff}}^2 = \sqrt{\frac{2k_{\text{B}}T_{\text{i}}}{M}} + v_{\text{nt}}^2 = \sqrt{v_{\text{th}}^2 + v_{\text{nt}}^2}$$

1. Assume  $T_{\text{i}}$  constant with height.

$$v_{\text{eff}}^2(R_1) - v_{\text{eff}}^2(R_2) = v_{\text{nt}}^2(R_1) - v_{\text{nt}}^2(R_2)$$

2. Assume waves undamped at low heights.

# Waves undamped at lowest heights



Undamped  $R < 1.15 R_{\odot}$  found by Banerjee et al. (1998, 2009), Dolla & Solomon (2008), Doyle et al. (1998)...

# Separate Thermal and Nonthermal

$$v_{\text{eff}}^2 = \sqrt{\frac{2k_{\text{B}}T_{\text{i}}}{M}} + v_{\text{nt}}^2 = \sqrt{v_{\text{th}}^2 + v_{\text{nt}}^2}$$

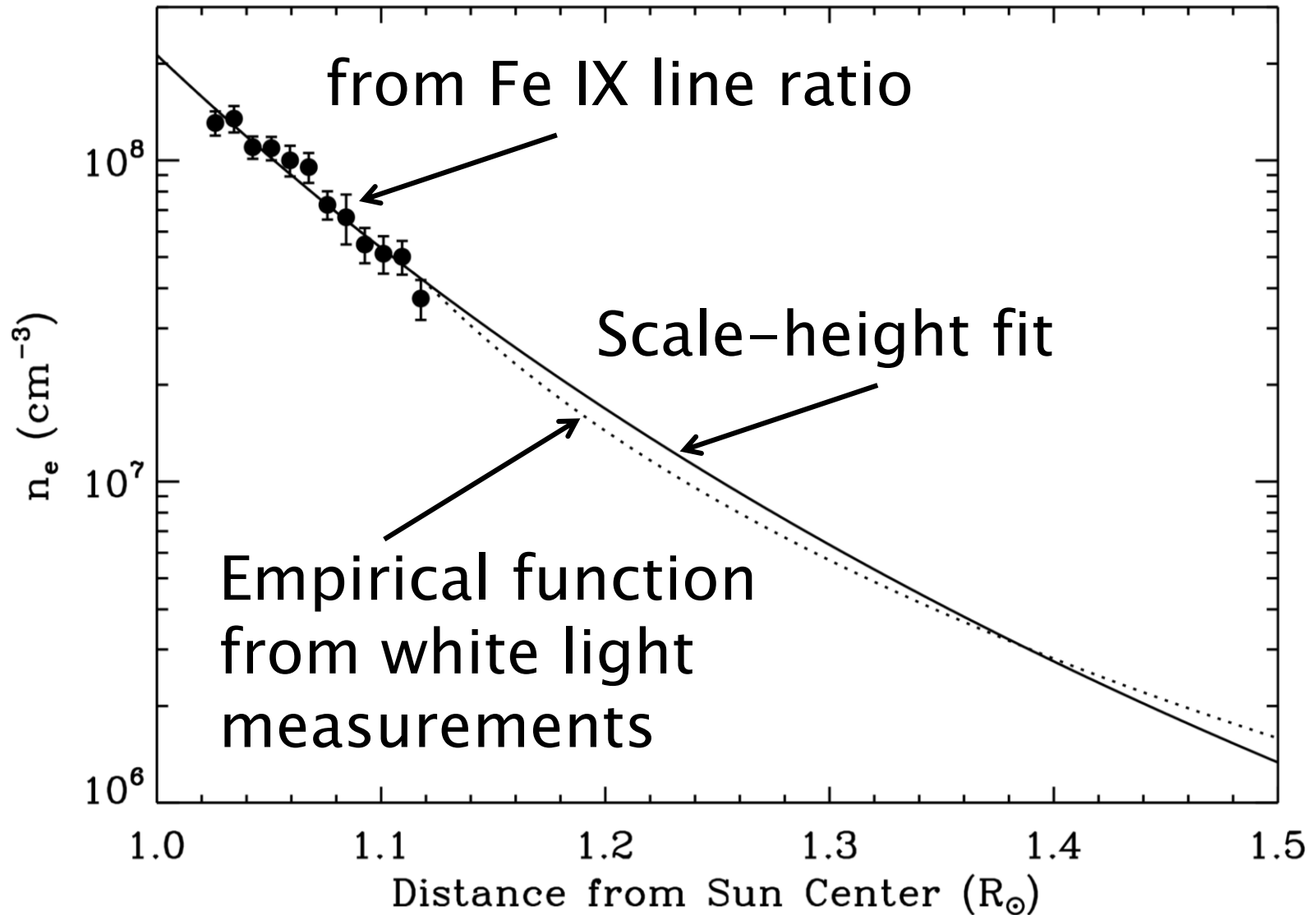
1. Assume  $T_{\text{i}}$  constant with height.

$$v_{\text{eff}}^2(R_1) - v_{\text{eff}}^2(R_2) = v_{\text{nt}}^2(R_1) - v_{\text{nt}}^2(R_2)$$

2. Assume waves undamped at low heights.

$$\frac{v_{\text{nt}}(R_1)}{v_{\text{nt}}(R_0)} = \left[ \frac{n_{\text{e}}(R_1)}{n_{\text{e}}(R_0)} \right]^{-1/4}$$

# Measured Density



# Separate Thermal and Nonthermal

$$v_{\text{eff}}^2 = \sqrt{\frac{2k_{\text{B}}T_{\text{i}}}{M}} + v_{\text{nt}}^2 = \sqrt{v_{\text{th}}^2 + v_{\text{nt}}^2}$$

1. Assume  $T_{\text{i}}$  constant with height.

$$v_{\text{eff}}^2(R_1) - v_{\text{eff}}^2(R_2) = v_{\text{nt}}^2(R_1) - v_{\text{nt}}^2(R_2)$$

2. Assume waves undamped at low heights.

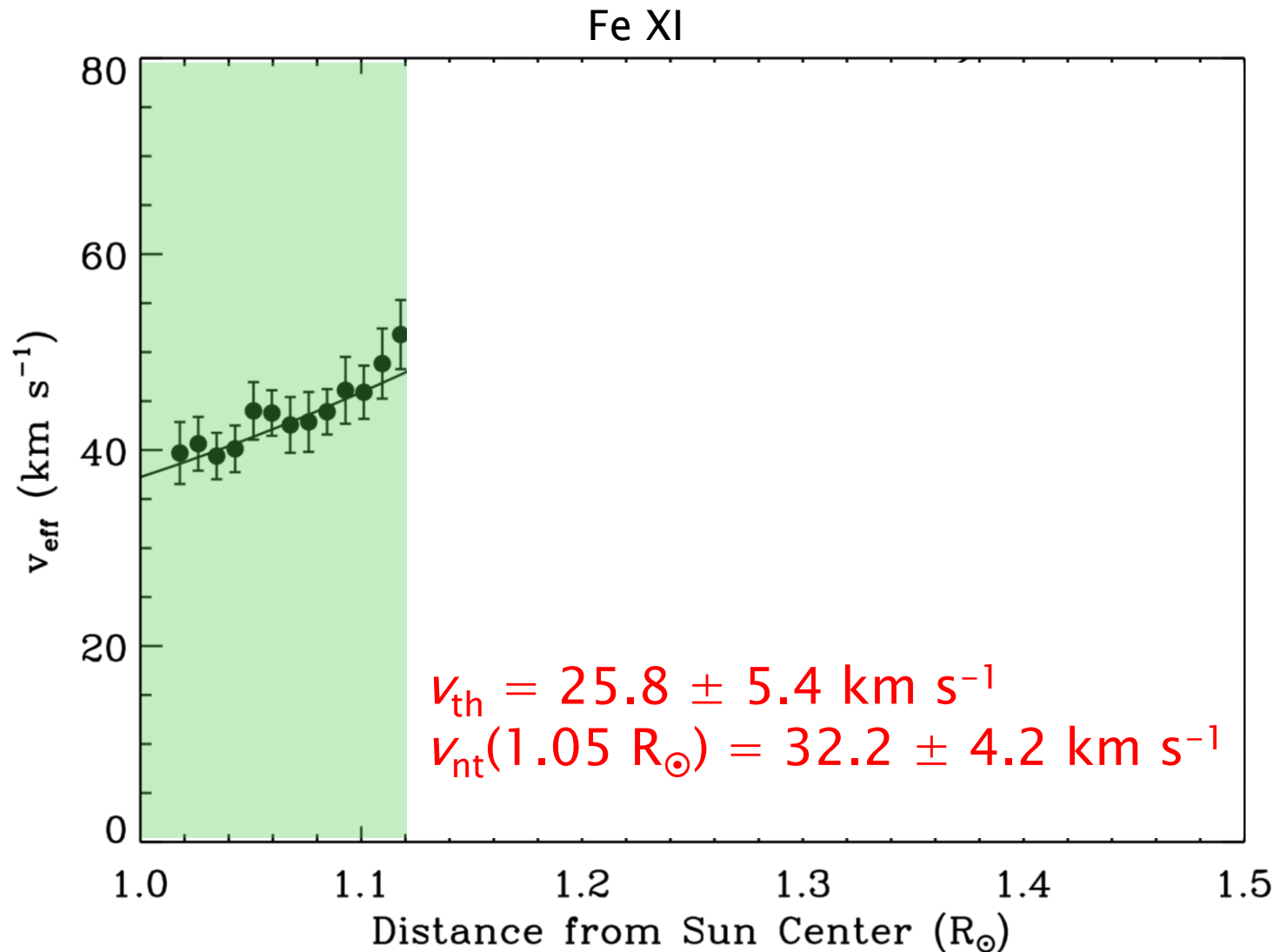
$$\frac{v_{\text{nt}}(R_1)}{v_{\text{nt}}(R_0)} = \left[ \frac{n_{\text{e}}(R_1)}{n_{\text{e}}(R_0)} \right]^{-1/4}$$

3. Fit data with  $v_{\text{th}}$  and  $v_{\text{nt}}(R_0)$  free parameters.

$$v_{\text{eff}}(R) = \sqrt{v_{\text{th}}^2 + v_{\text{nt}}^2(R_0) \left[ \frac{n_{\text{e}}(R)}{n_{\text{e}}(R_0)} \right]^{-1/2}}$$



# Example Fit



# Extend to larger heights

The fit used up to  $1.12 R_{\odot}$ .

To extend, assume  $T_i$  remains constant.

In reality,  $T_i$  increases at large heights.

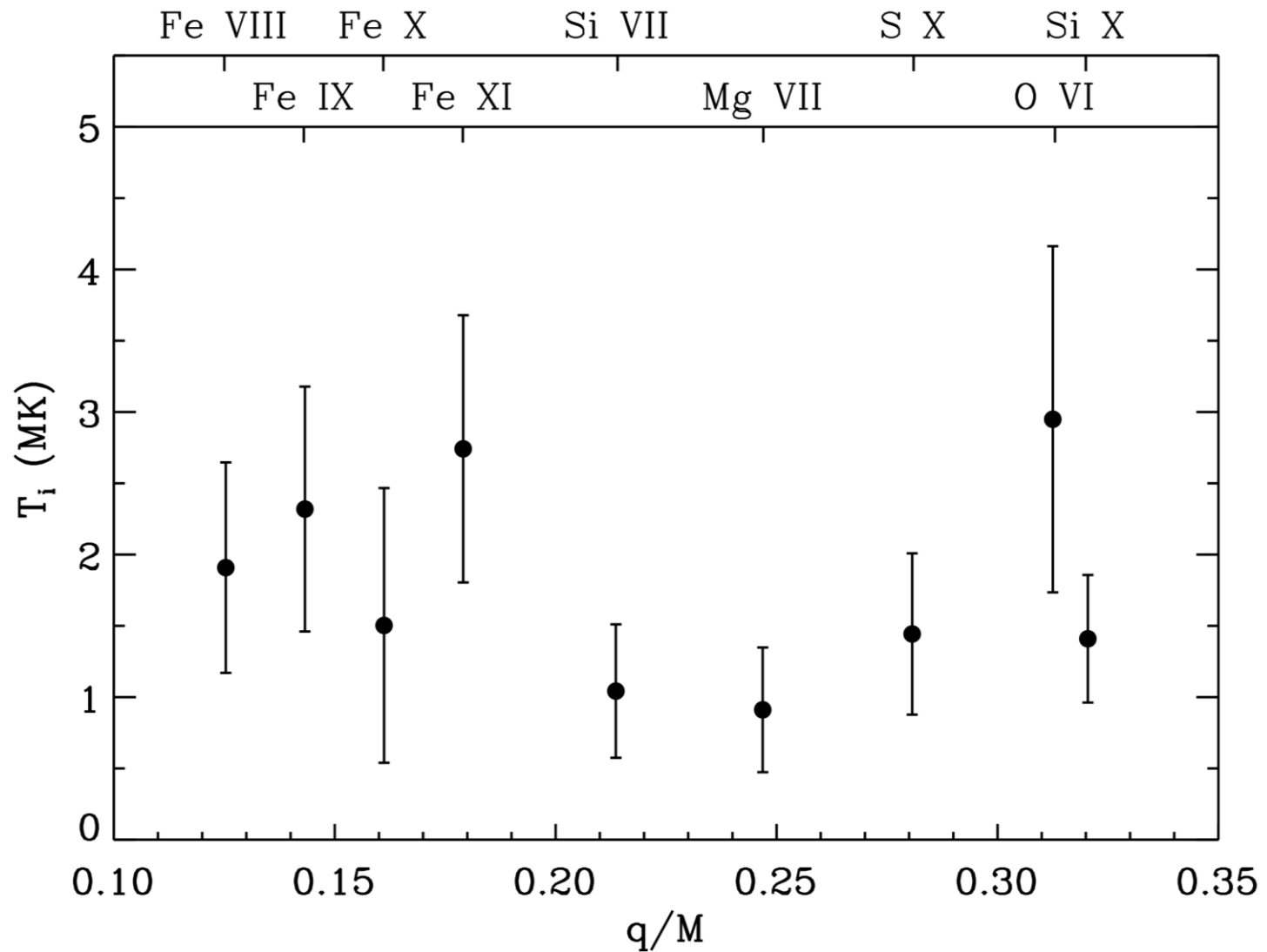
In deriving  $v_{nt}$ , not subtracting enough  $v_{th}$ .

We may underestimate the actual damping.

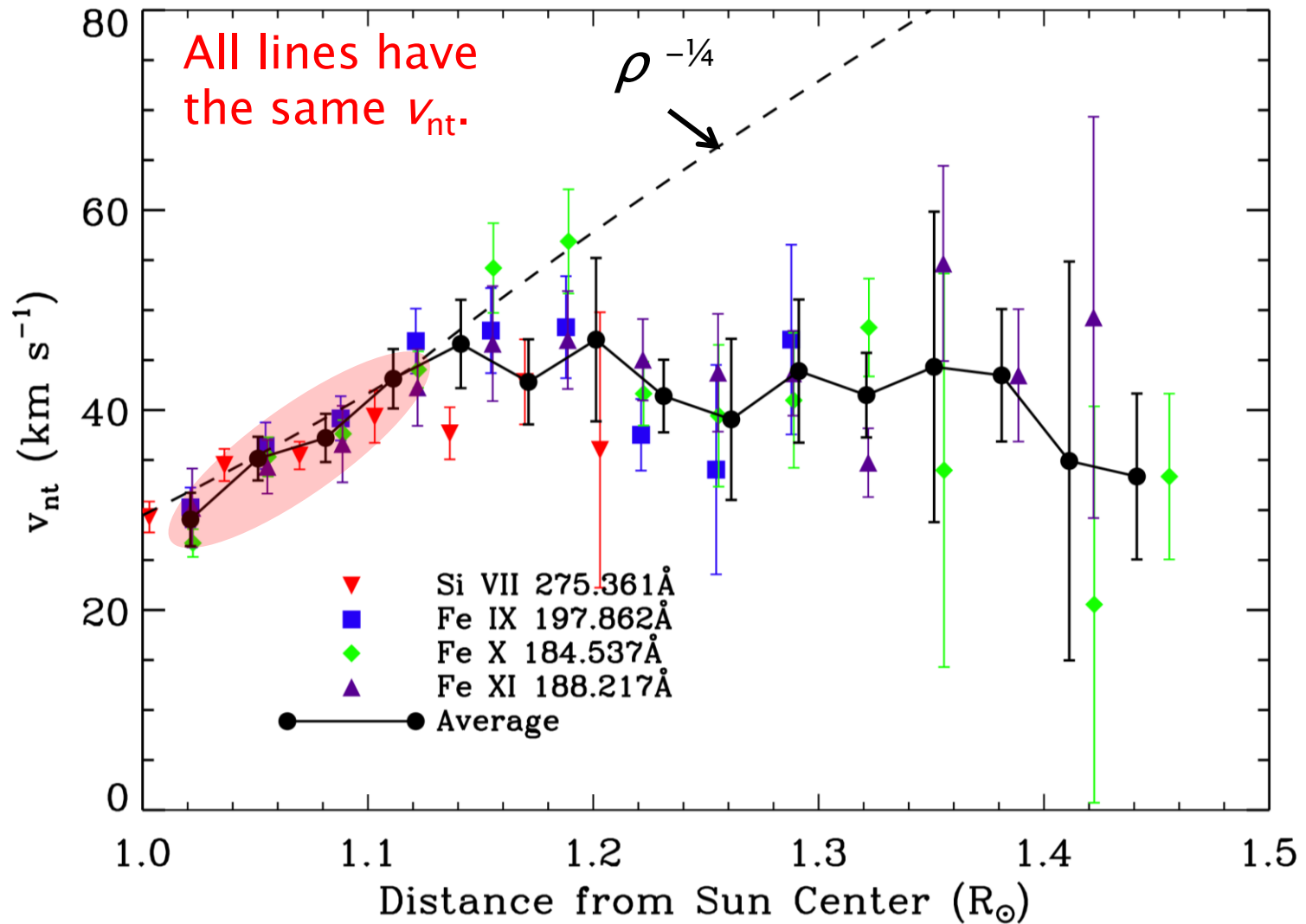
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# Ion Temperature



# Damping begins about 1.15 $R_{\odot}$



# Wave Energy Flux

$$F = 2\rho v_{nt}^2 V_A$$

$$V_A = \frac{B}{\sqrt{4\pi\rho}}$$

$$\rho \sim 1.15 m_p n_e$$

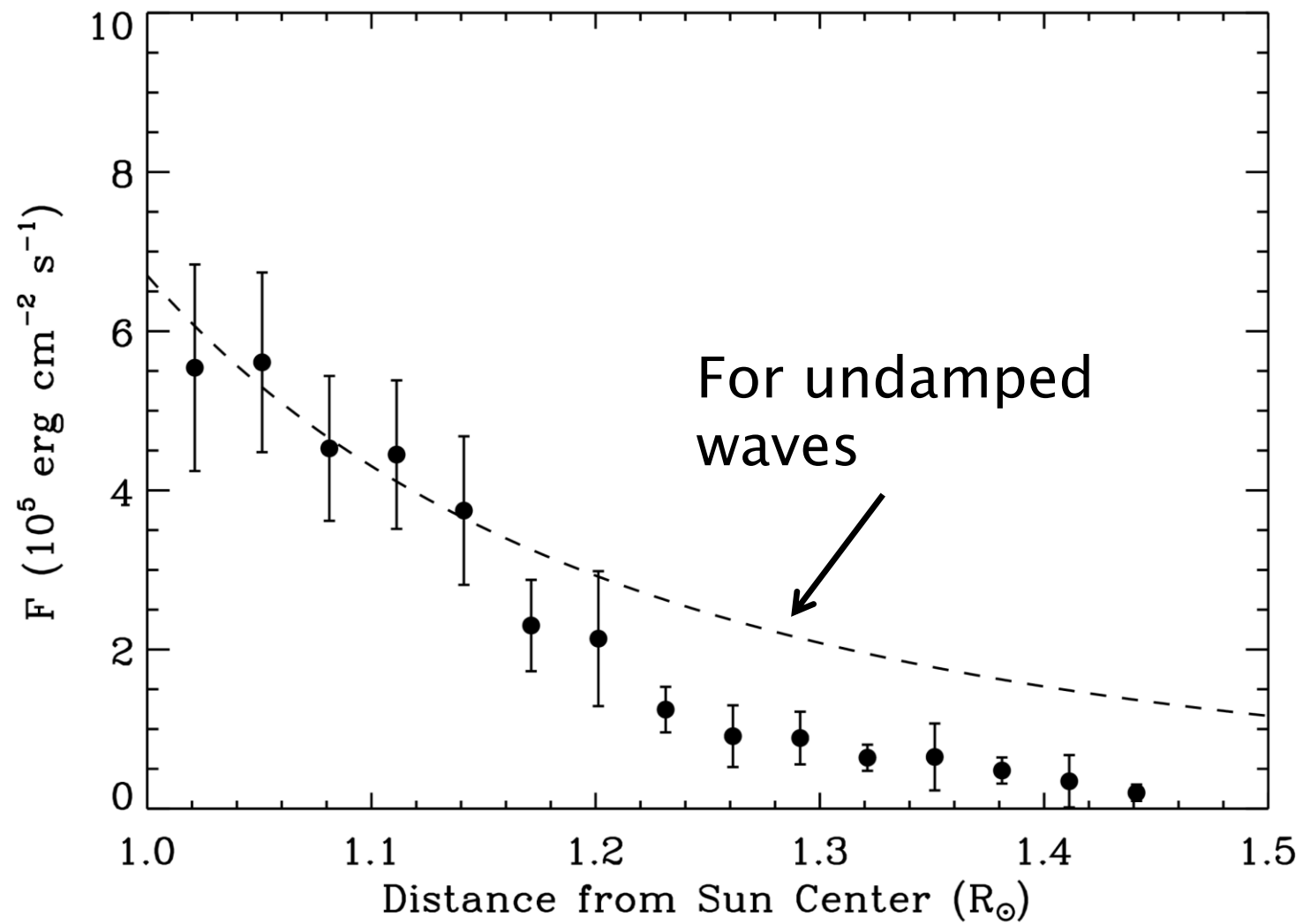
For  $B$  use empirical model (Cranmer et al. 1999) .

$$B(R)/B(R_\odot) = A(R_\odot)/A(R)$$

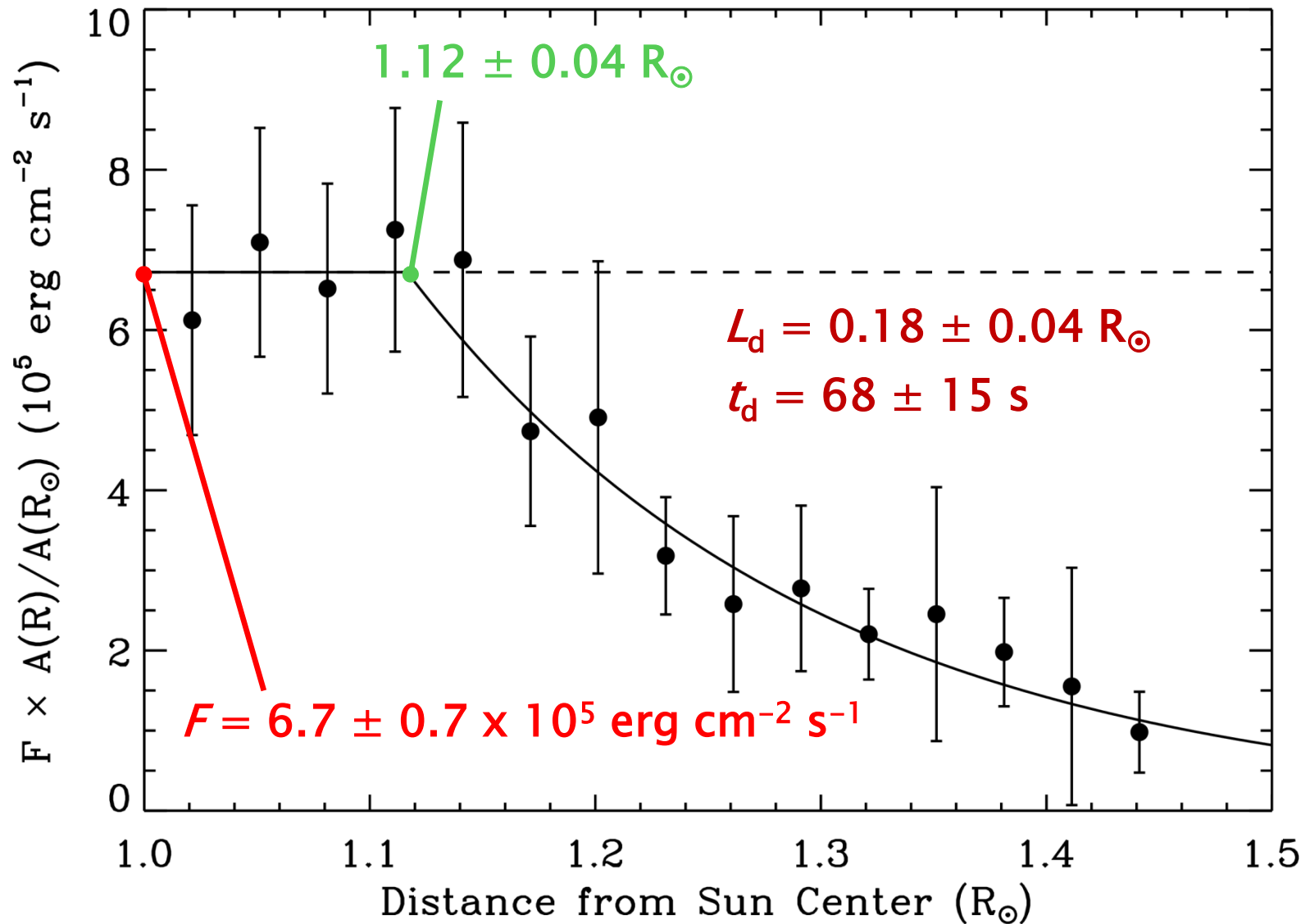
$$B(R_\odot) = 7.3 \pm 1 \text{ G (Wang 2010)}$$



# Wave Energy Flux $F$



# $F$ Corrected for Area Expansion



# Energy Required

To heat coronal hole and accelerate solar wind  
 $\sim 6 \times 10^5 \text{ erg cm}^{-2} \text{ s}^{-1}$  required during this time.

Measured  $6.7 \pm 0.7 \times 10^5 \text{ erg cm}^{-2} \text{ s}^{-1}$

Waves carry and dissipate enough energy at low enough heights to heat coronal holes and accelerate the fast solar wind.

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# Damping Mechanisms

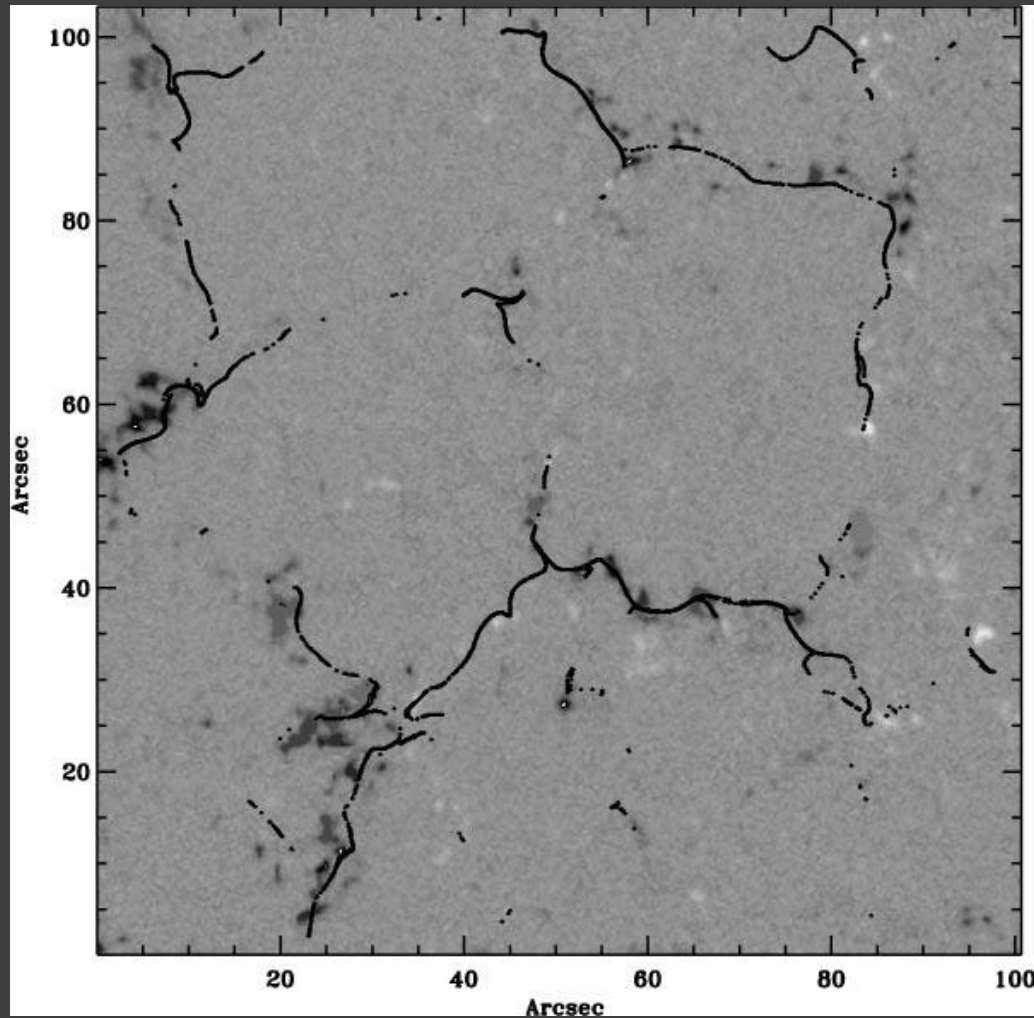
- ▣ The structure of the real solar magnetic field is not homogeneous – allows additional damping mechanisms:
- ▣ Gradients across the magnetic field:
  - E.g. Phase mixing
- ▣ Gradients along the field:
  - E.g. Wave reflection & turbulence

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# Structure of Surface Magnetic Field



Roudier et al. (2009)

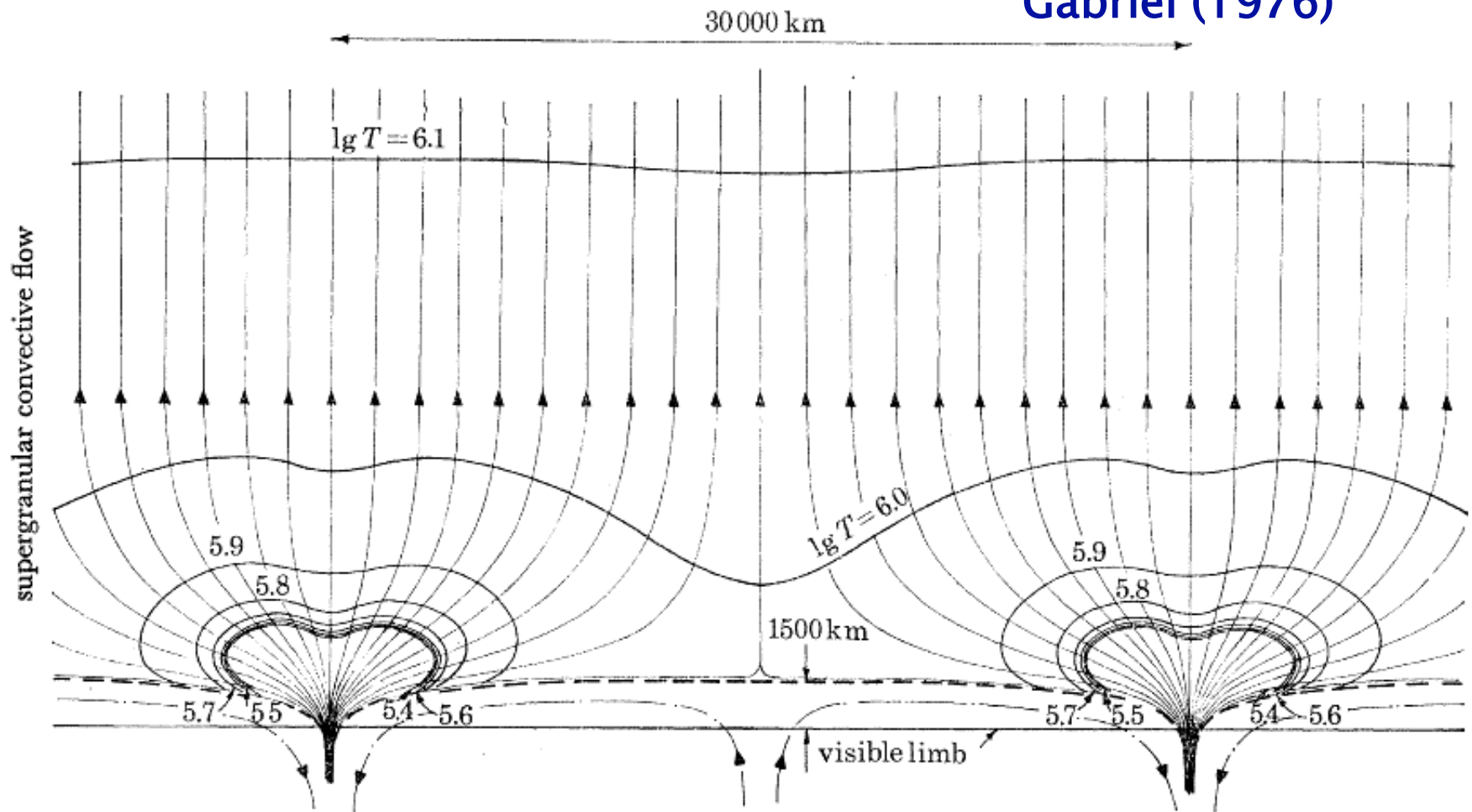
Greyscale is magnetic field, white = positive, black = negative, grey = neutral.

Black lines outline edges of convection cells (supergranules).

Magnetic network.

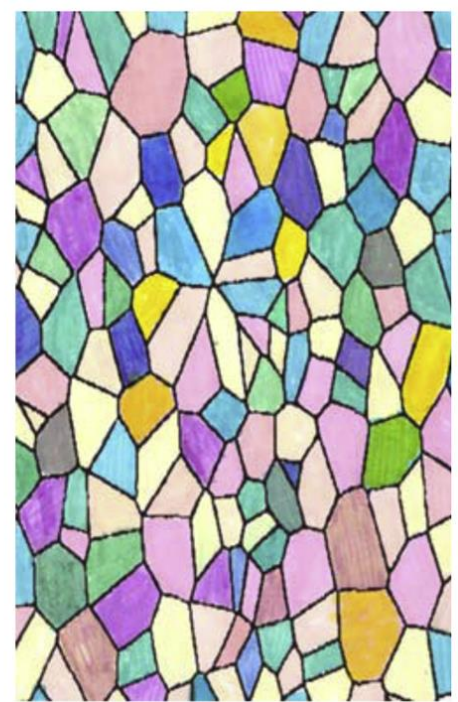
# Structure of Solar Magnetic Field

Gabriel (1976)



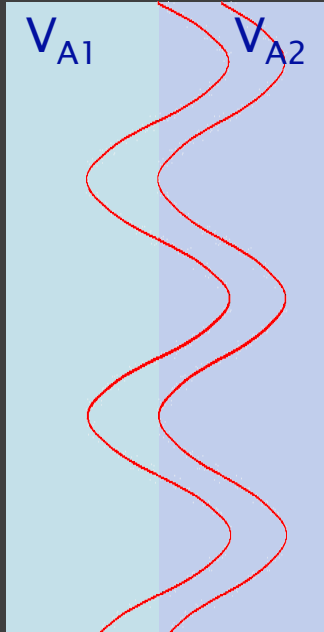
# Structure of Solar Magnetic Field

Sketch of flux tubes at 1 AU

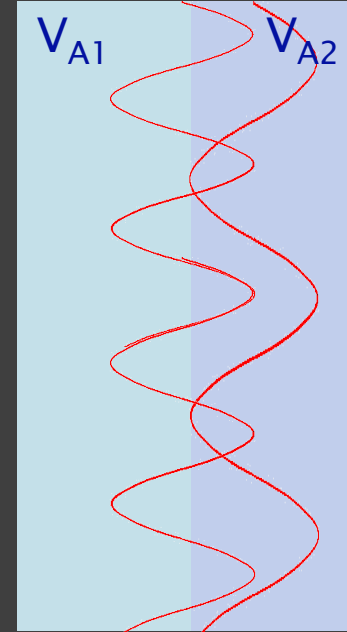


Borovsky (2008)

# Phase Mixing



Footpoint oscillations generate waves at same frequency, initially in-phase.



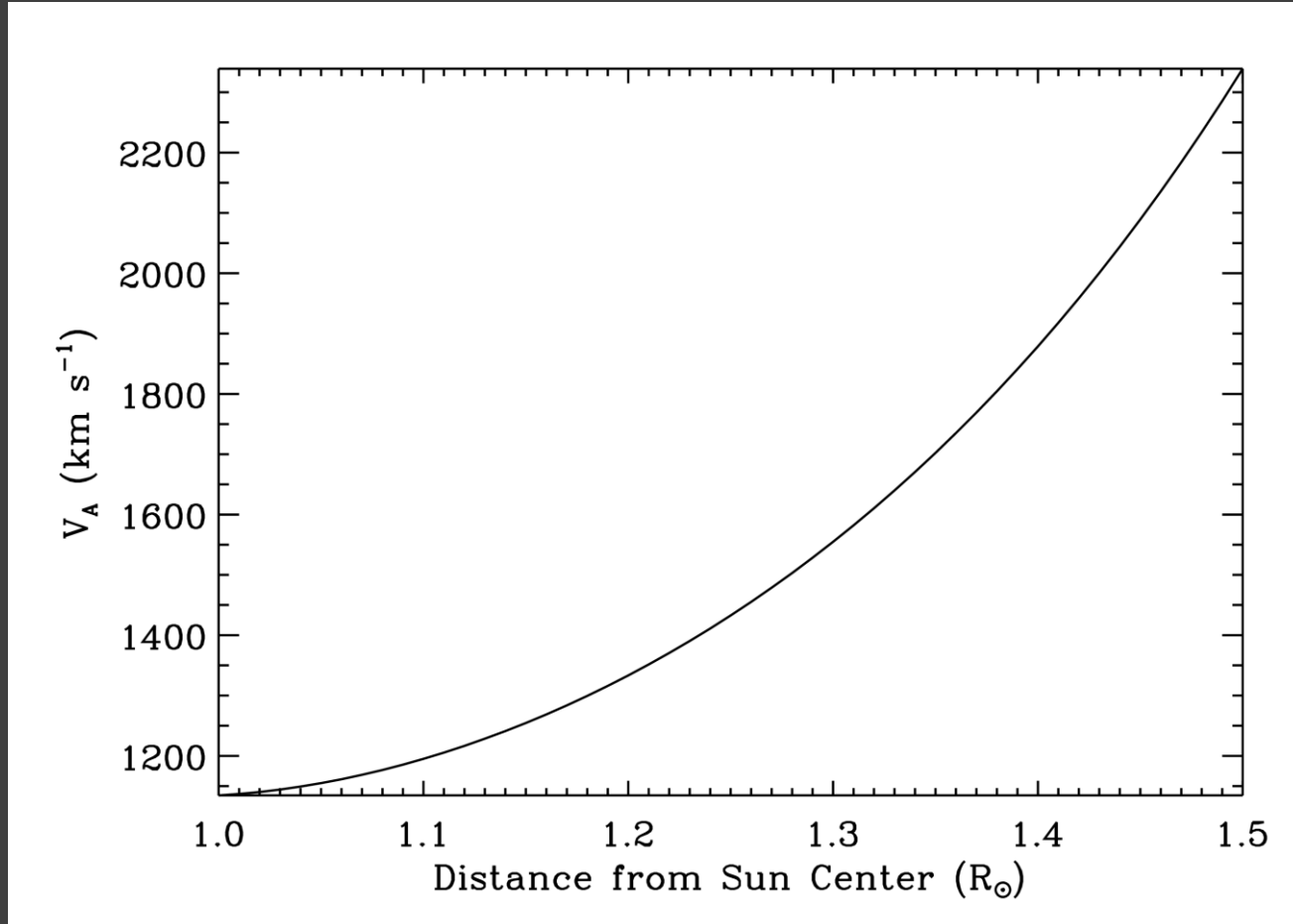
Due to different  $V_A$  (phase velocity) waves become out of phase.

Increases viscous damping (friction)

# Outline

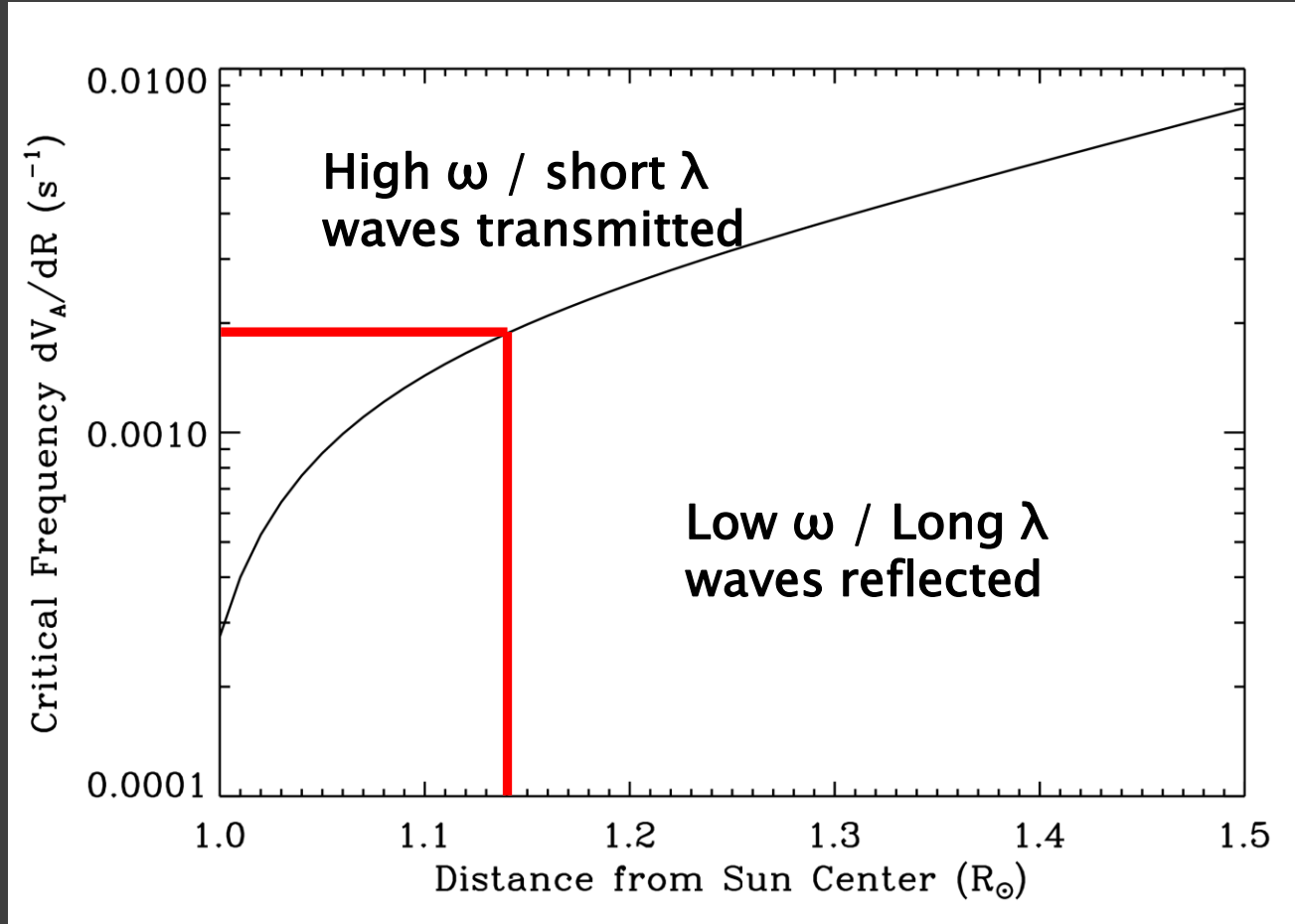
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# Wave Reflection and Turbulence



Strong wave reflection if  $V_A$  changes over scales shorter than the wavelength

# Wave Reflection & Turbulence



Waves reflected if  $\omega < \omega_{\text{critical}} \sim dV_A/dr$

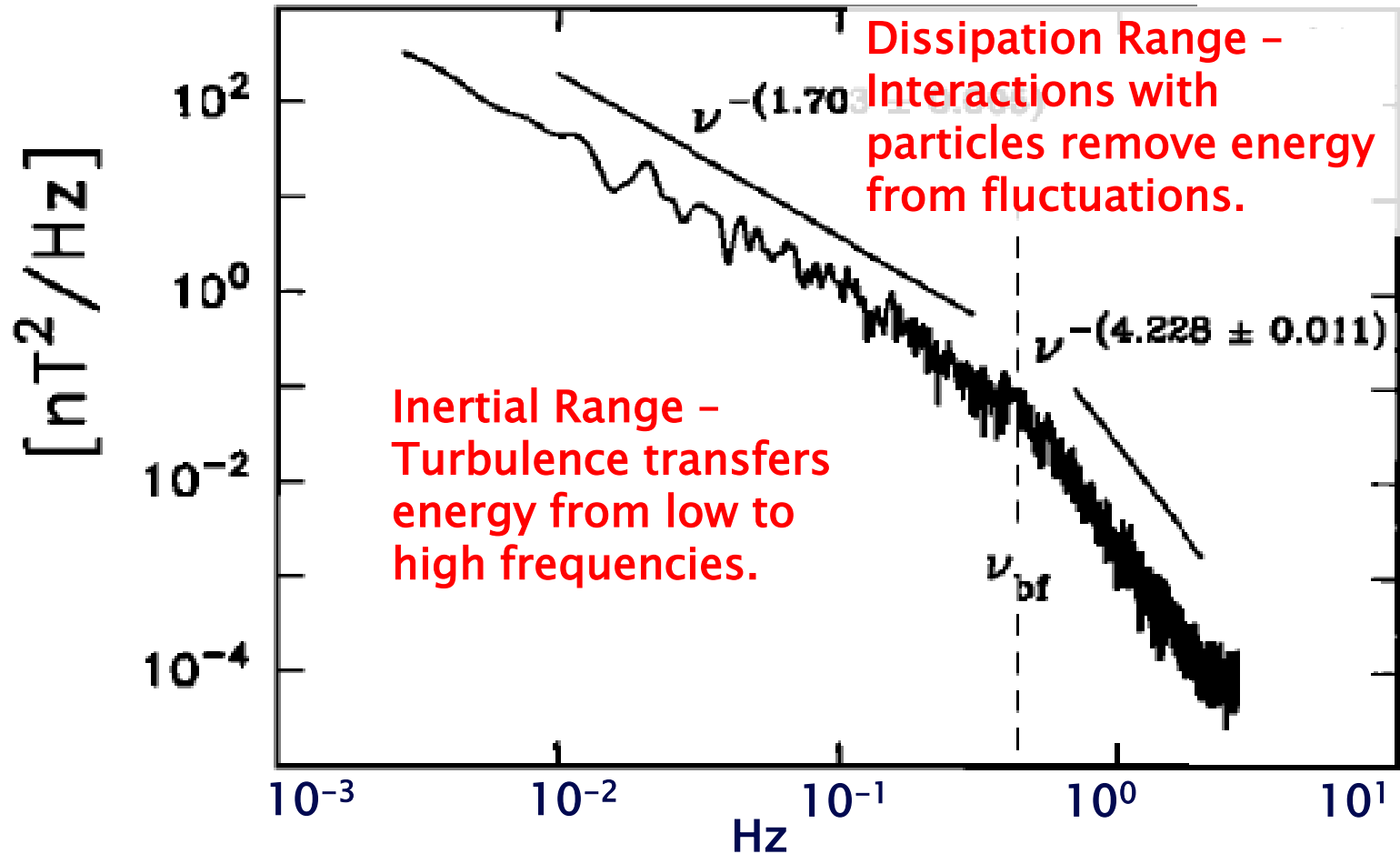
Photospheric source  $\omega \sim 0.05\text{--}0.005 s^{-1}$



# Turbulence in the Solar Wind

Leamon et al. (1998) – Measurements at 1 AU

Power Density



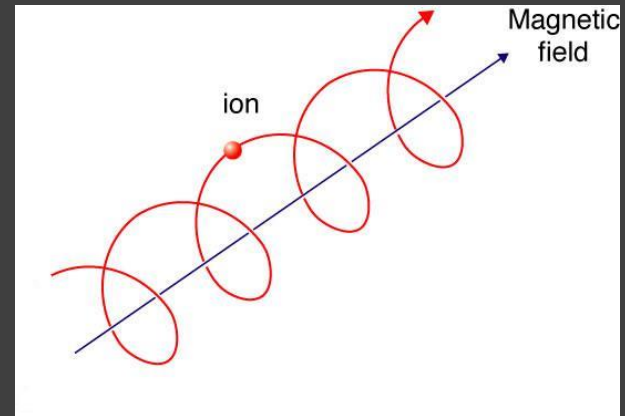
Interaction of the outward propagating and reflected waves generates turbulence.



# Ion Cyclotron Resonant Heating

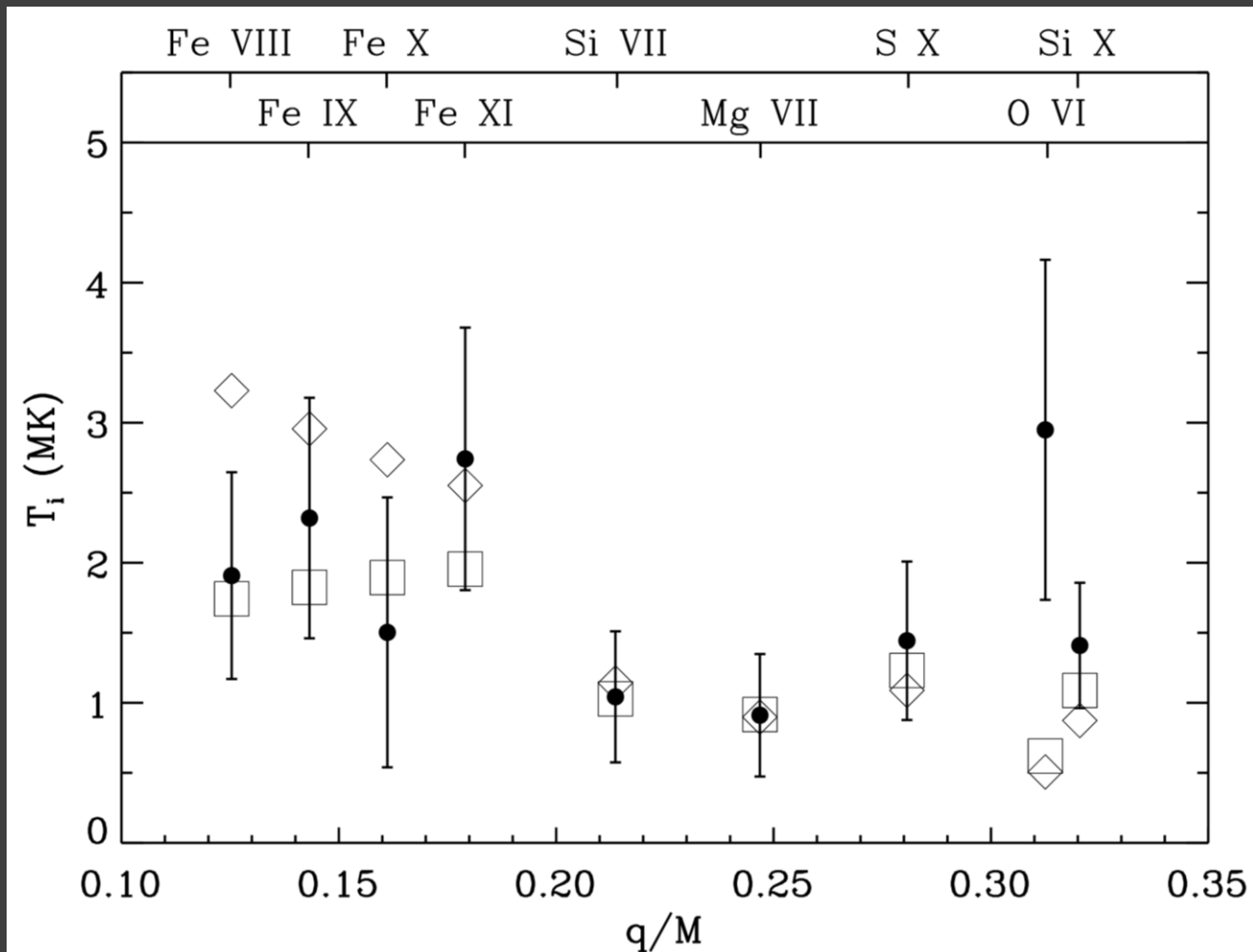
- Ion cyclotron frequency

$$\Omega_i = qB/M.$$



- If wave frequency close to  $\Omega_i$  then waves can efficiently transfer energy to the ions.
- Expect more power at low frequencies, so heating depends on  $q/M$ .
- Measure  $T_i$  vs.  $q/M$  to look for evidence of resonant heating.

# Ion Temperature



Measured  $T_i$  compared to predicted trends for turbulence-driven heating

# Conclusions and Open Questions

- ▣ Waves carry and dissipate enough energy at low enough heights to heat coronal holes and accelerate the fast solar wind.
- ▣ What causes the wave damping?
  - Magnetic structure allows damping processes.
  - Laboratory experiments are planned to study these.
- ▣ What heats other structures of the corona?
  - Models and some measurements suggest quiet Sun could also be heated by waves.
  - Currently applying this type of analysis to quiet Sun.